

Rigorous computations for infinite dimensional dynamical systems

Jean-Philippe Lessard



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April 12th, 2013*

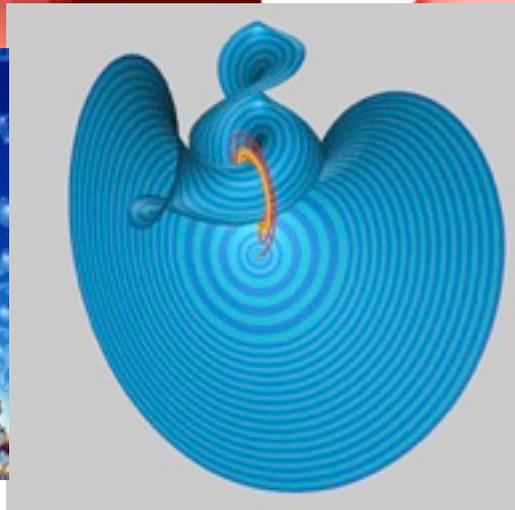
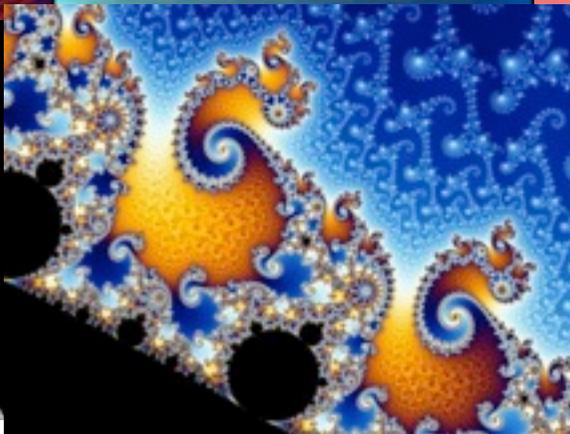
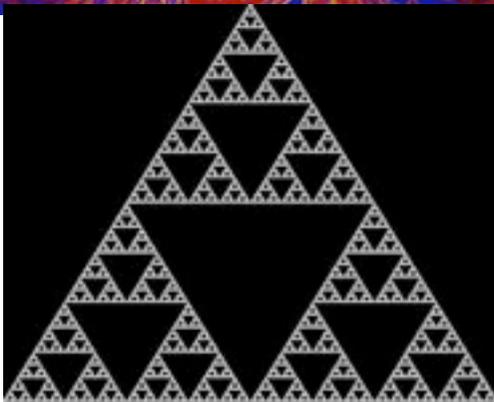
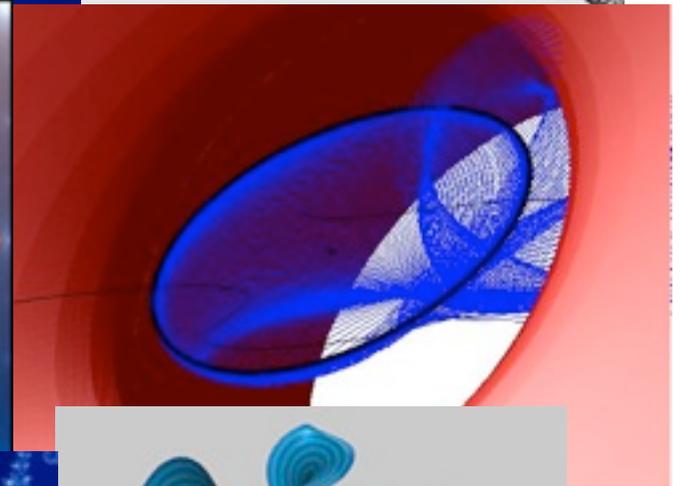
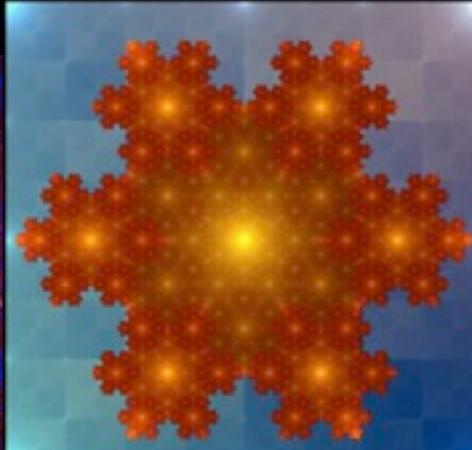
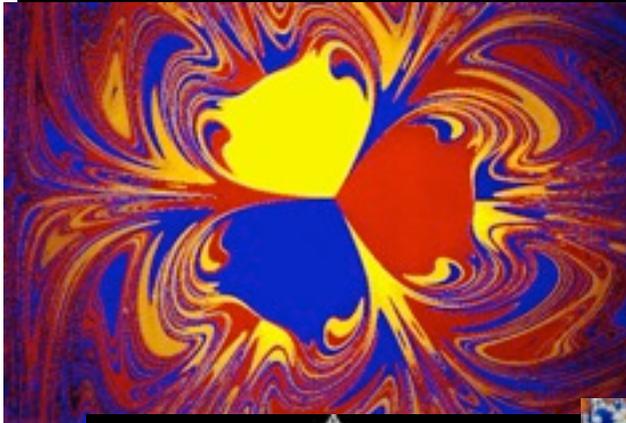
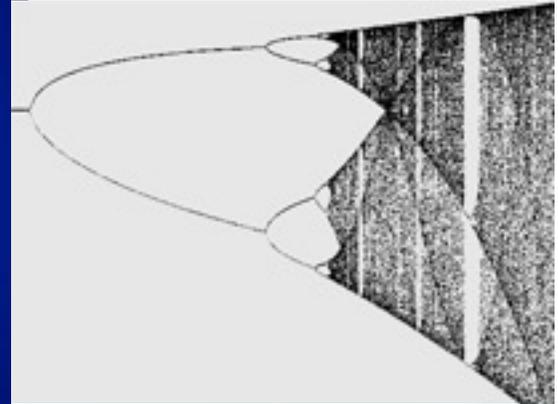
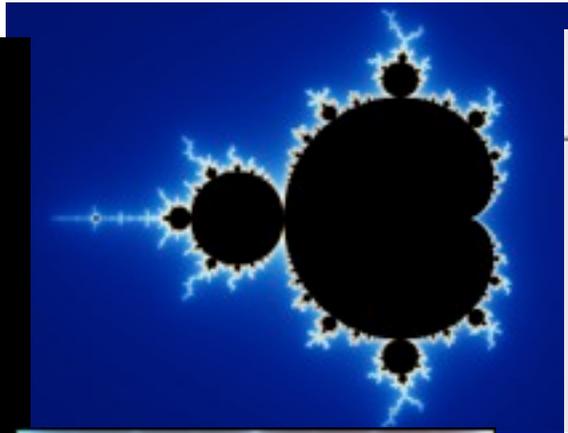
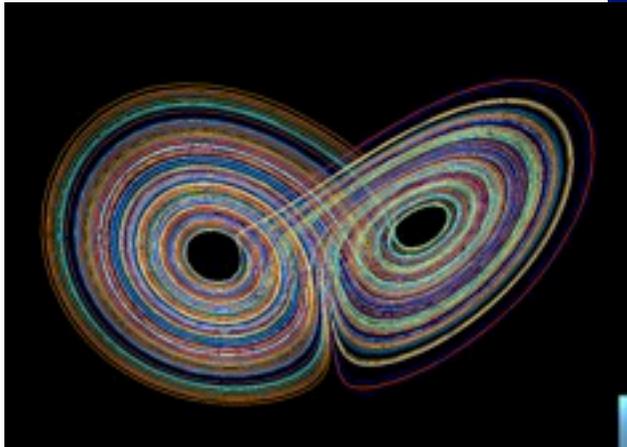
What is a dynamical system?

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Popular answer: a math subject that produces beautiful pictures!

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Informal answer: a system that evolves with time

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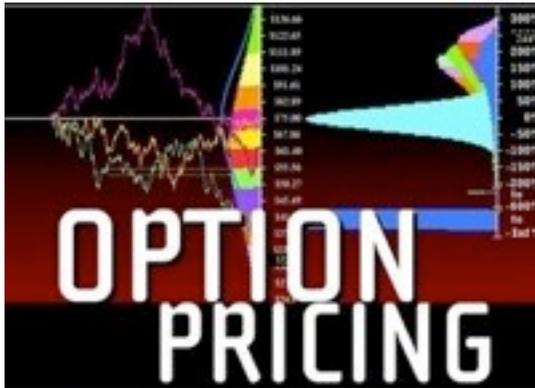
The dynamical system concept is a mathematical formalization for any fixed "rule" which describes the time dependence of a point's position in its ambient space.

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Informal answer: a system that evolves with time



The dynamical system concept is a mathematical formalization for any fixed "rule" which describes the time dependence of a point's position in its ambient space.



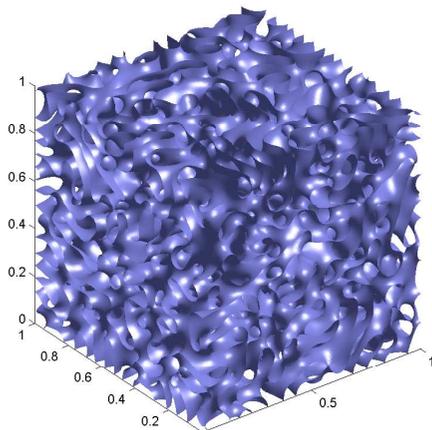
finance



fluids



weather prediction



material science



population dynamics



chemical reactions

What is a dynamical system?

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Formal answer: a math definition

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A dynamical system is a tuple (T, M, Φ)

T : monoid (**time**)

M : set (**state space**)

Φ : map (**evolution function**)

$$\Phi : T \times M \rightarrow M$$

satisfying the two following properties

$$\begin{cases} \Phi(0, x) = x \\ \Phi(t_2, \Phi(t_1, x)) = \Phi(t_1 + t_2, x) \end{cases}$$

$$\forall x \in M \text{ and } \forall t_1, t_2 \in T$$

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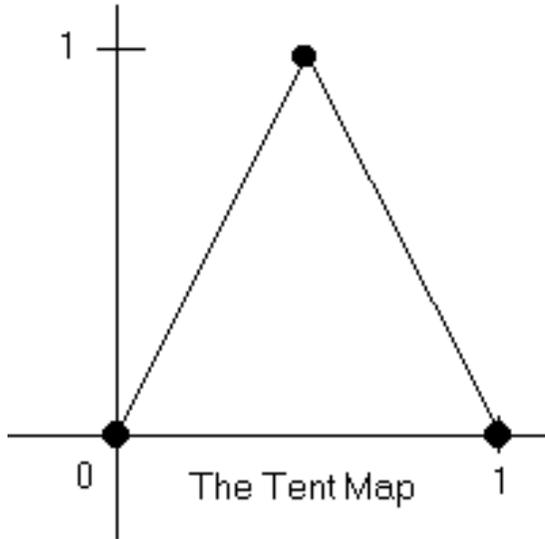
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In case the state space M is a function space, we have an infinite dimensional dynamical system !

Examples

I. Finite dimensional discrete dynamical systems



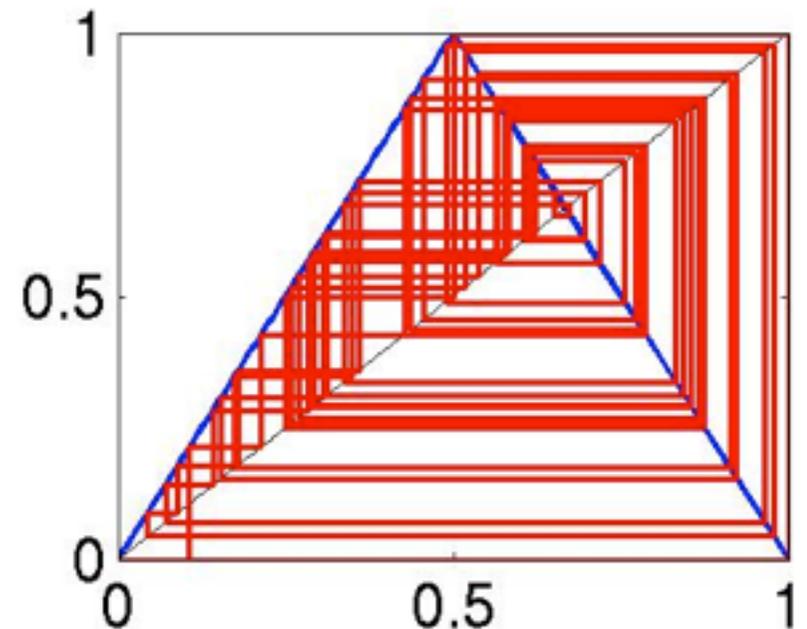
$$f(x) = \begin{cases} 2x, & \text{for } x \in [0, \frac{1}{2}) \\ 2(1-x), & \text{for } x \in [\frac{1}{2}, 1] \end{cases}$$

$T = \mathbb{N}$ (discrete time)

$M = [0, 1]$ (state space)

$\Phi : T \times M \rightarrow M$

$(n, x) \mapsto \Phi(n, x) = f^n(x)$



Examples

2. Finite dimensional continuous dynamical systems: ODEs

$$\text{(IVP)} \begin{cases} \frac{dx}{dt} = f(x) \\ x(0) = x_0 \end{cases} \quad [f \in C^1(\mathbb{R}^n)]$$

$\Phi(t, x_0)$: solution of the (IVP)

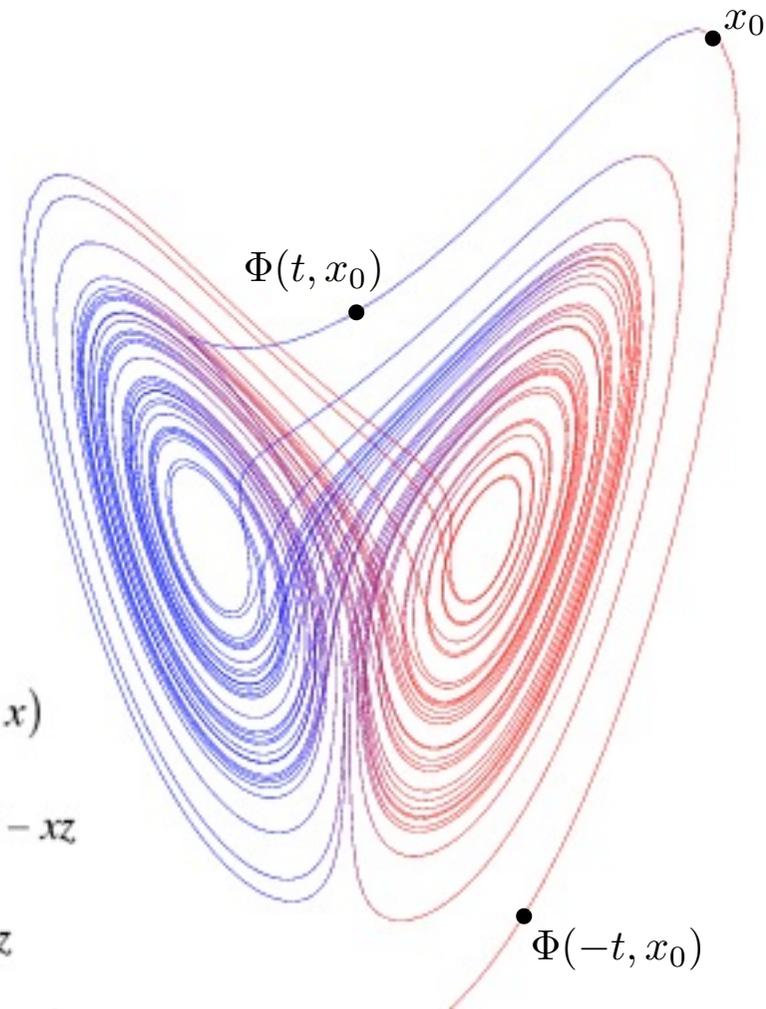
$T = \mathbb{R}$ (continuous time)

$M = \mathbb{R}^n$ (state space)

$$\begin{aligned} \Phi : T \times M &\rightarrow M \\ (t, x_0) &\mapsto \Phi(t, x_0) \end{aligned}$$

$$\begin{aligned} \frac{dx}{dt} &= \sigma(y - x) \\ \frac{dy}{dt} &= rx - y - xz \\ \frac{dz}{dt} &= xy - bz \end{aligned}$$

Lorenz equations



Examples

3. Infinite dimensional continuous dynamical systems

(a) Partial differential equations

Cahn-Hilliard equation

$$\frac{\partial u}{\partial t} - \Delta \left(-\nu \Delta u - u + u^3 \right) = 0$$

$$\Omega \subset \mathbb{R}^n, \quad n = 1, 2, 3$$

$$T = [0, \infty) \text{ (continuous time)}$$

$$M = L^2(\Omega) \text{ (infinite dimensional state space)}$$

$$\Phi : T \times M \rightarrow M$$

$$(t, u_0) \mapsto \Phi(t, u_0) \text{ (semigroup)}$$

Examples

3. Infinite dimensional continuous dynamical systems

(b) Delay differential equations $y'(t) = \mathcal{F}(y(t), y(t - \tau))$

$T = [0, \infty)$ (continuous time)

$M = C[-\tau, 0]$ (infinite dimensional state space)

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Examples

3. Infinite dimensional continuous dynamical systems

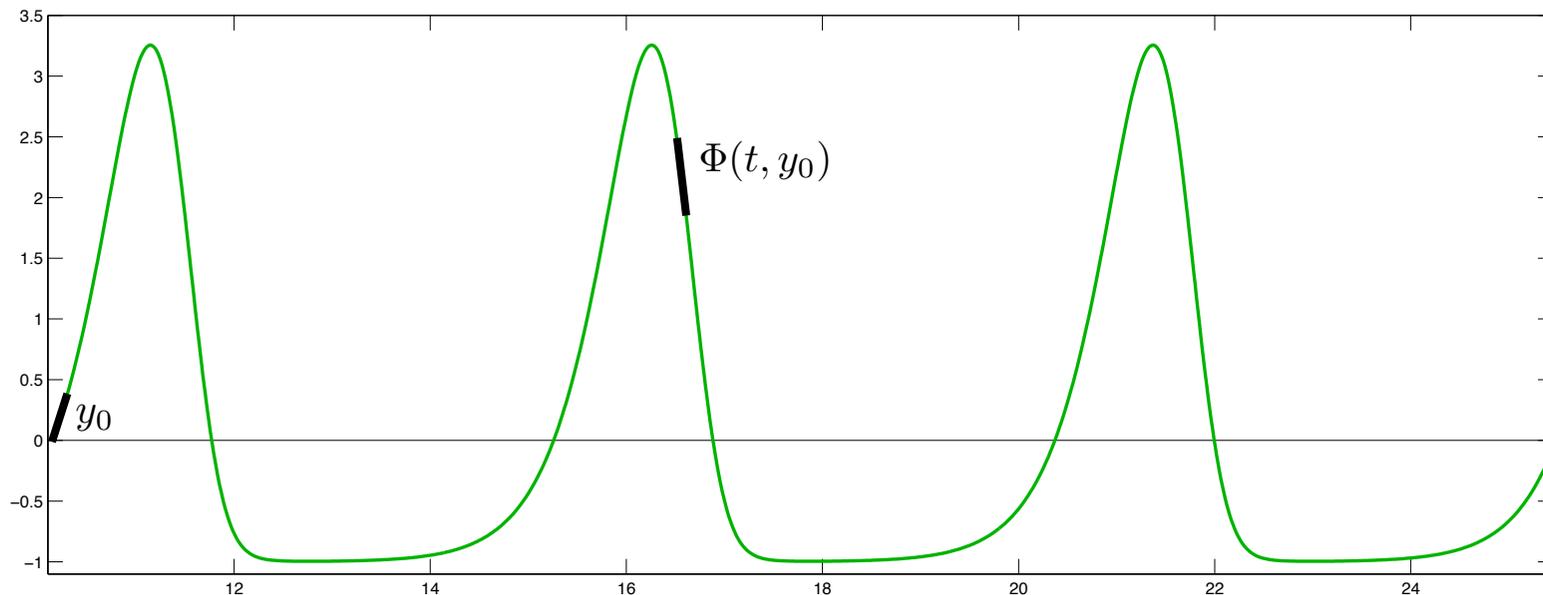
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ex: $y'(t) = -\frac{12}{5}y(t-1)[1+y(t)]$

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In any dynamical system, it is the bounded solutions which are most important and which should be investigated first.



Henri Poincaré

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Compact invariant sets

Exploit smoothness, boundedness and low dimensionality.

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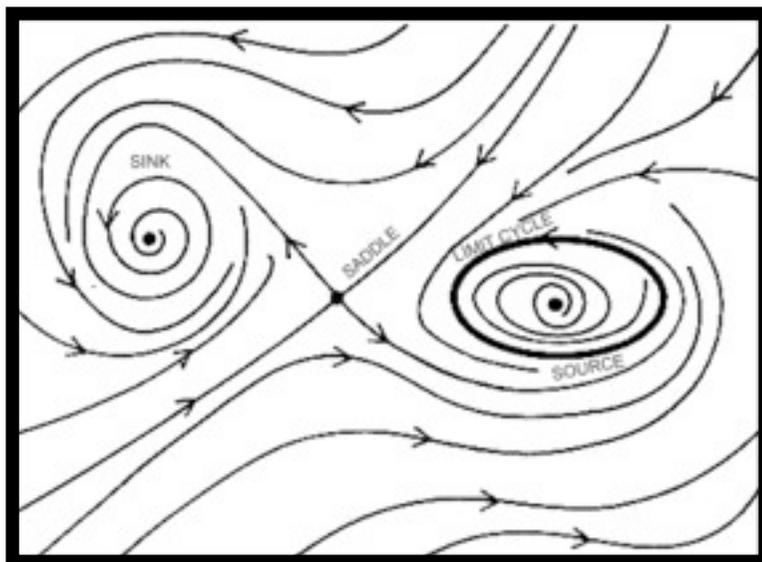
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- Equilibrium solutions.
- Time periodic solutions.
- Connecting orbits.
- Global attractors.

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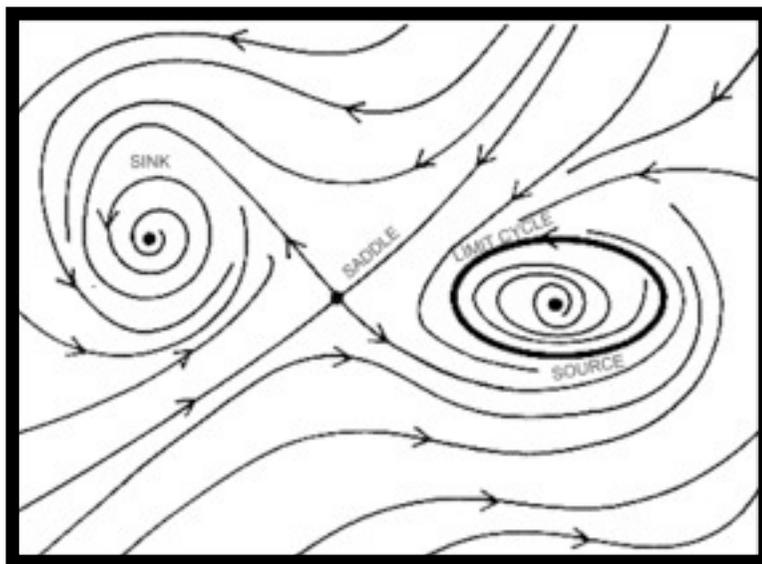
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A standard approach is to get insight from numerical simulations to formulate new conjectures, and then attempt to prove the conjectures using pure mathematical techniques only. Actually, this strong dichotomy need not exist in the context of dynamical systems, as the strength of numerical analysis and functional analysis can be combined to prove, in a **rigorous mathematical sense**, the existence of equilibria, periodic solutions, connecting orbits... and even chaotic dynamics !

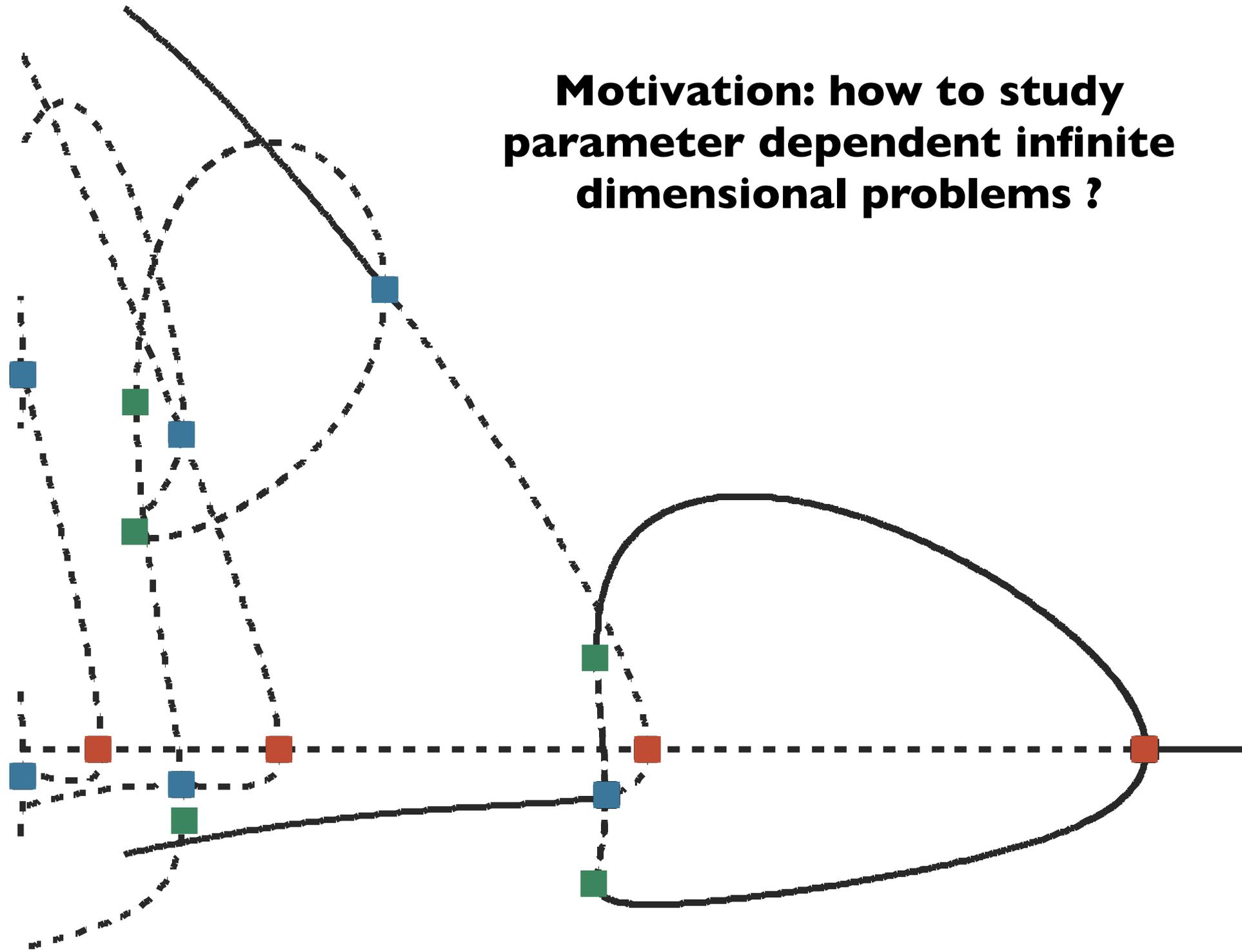
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Rigorous computations

The goal of rigorous computations is to construct algorithms that provide an approximate solution to the problem together with precise and possibly efficient bounds within which the exact solution is guaranteed to exist in the mathematically rigorous sense.

**Motivation: how to study
parameter dependent infinite
dimensional problems ?**



$$\mathcal{F}(x) = 0$$

 X

● x_1

● x_3

● x_2

● x_4

● x_6

● x_5

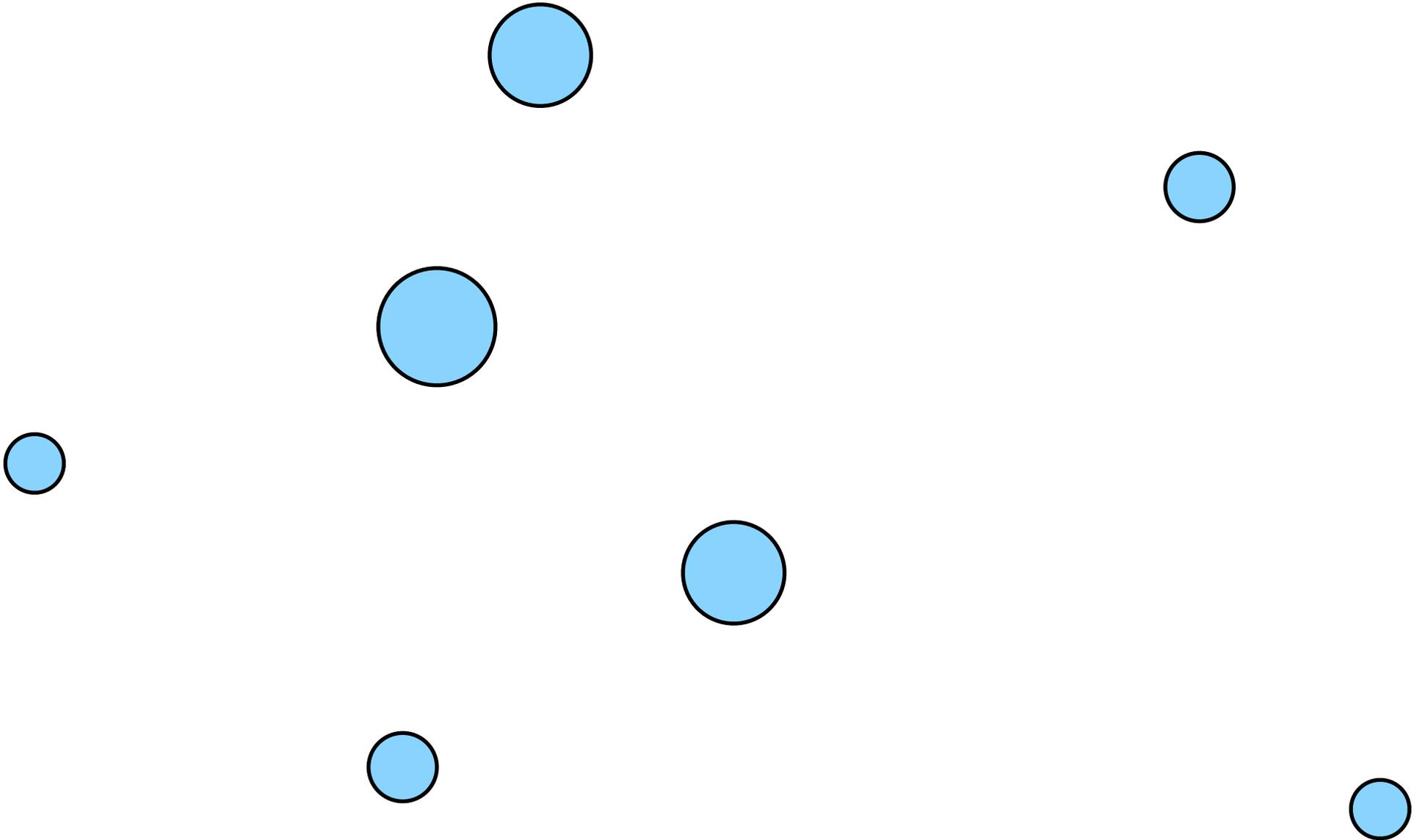
● x_7

$$\mathcal{F}(x) = 0$$

 X x_1 x_3 x_2 x_4 x_6 x_5 x_7

Often impossible to compute exactly !

$$\mathcal{F}(x) = 0$$

 X 

Alternative: find small balls in which it is demonstrated (in a mathematically rigorous sense) that a unique solution exists.

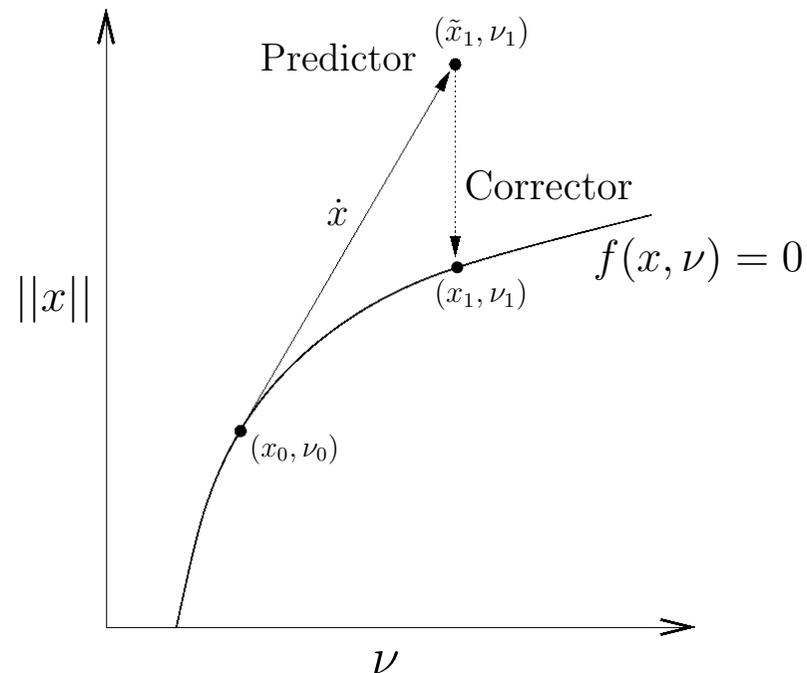
Rigorous Computations (Ingredients)

1. Smoothness of the solutions
2. Banach space of algebraically decaying sequences
3. Finite dimensional Galerkin projection
4. Bounds on the truncation error terms (Analytic estimates)
5. Fixed point theory, Uniform contraction principle
6. Numerical analysis (continuation, Fast Fourier transform)
7. Interval Arithmetic

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Continuation (Predictor-Corrector Algorithm)



Rigorous Computations

$$F(u, \nu) = 0$$

(Differential Equation)

spectral method



$$f(x, \nu) = 0$$

x : modes

ν : parameter

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Knowledge
about regularity



$$x \in \Omega^s = \left\{ (x_k)_k : \|x\|_s = \sup_k \{ \|x\|_\infty k^s \} < \infty \right\}$$

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Consider \bar{x} such that $f^{(m)}(\bar{x}, \nu_0) \approx 0$.

Galerkin approximation

$$f(x, \nu) = 0 \iff T_\nu(x) = x$$

Newton-like operator at \bar{x}

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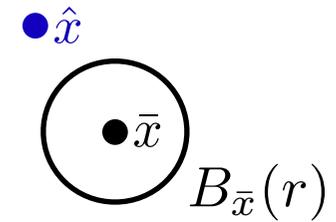
$$T_\nu : \Omega^s \rightarrow \Omega^s$$

$$T_\nu(x) = x - Jf(x, \nu)$$

$$J \approx D_x f(\bar{x}, \nu_0)^{-1}$$

The chances of contracting a small set B around \bar{x} depends on the magnitude of the eigenvalues of J .

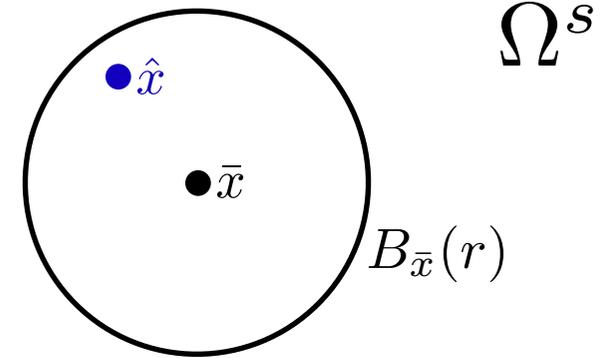
Q: How to find a ball $B_{\bar{x}}(r)$ such that
 $T_\nu: B_{\bar{x}}(r) \rightarrow B_{\bar{x}}(r)$ is a contraction?



Ω^s

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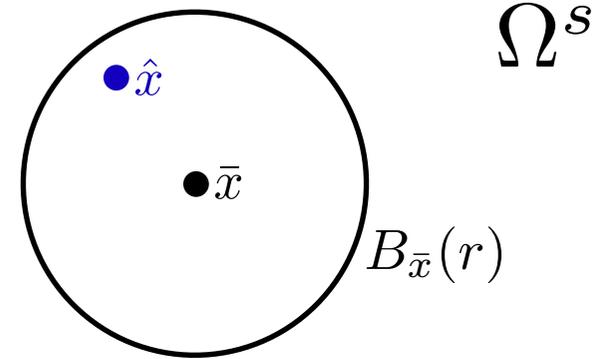
$$B_{\bar{x}}(r) = \bar{x} + B(r)$$



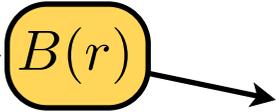
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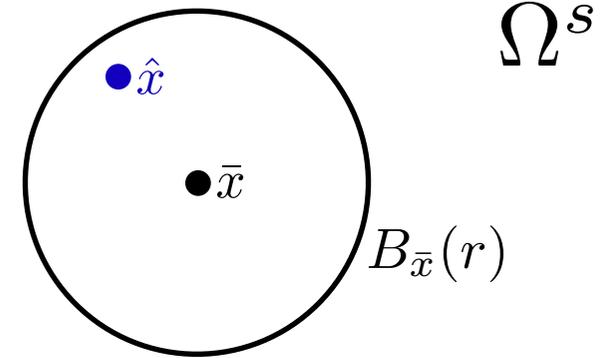
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Ball of radius r centered at 0 in the space Ω^s



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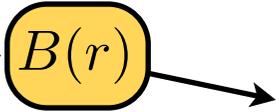
$B_{\bar{x}}(r) = \bar{x} + B(r)$  Ball of radius r centered at 0 in the space Ω^s

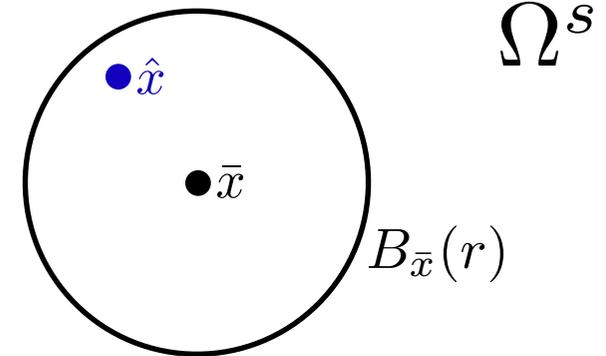


A: **Radii polynomials** $\{p_k(r)\}_k$: upper bounds satisfying

$$\left| [T_\nu(\bar{x}) - \bar{x}]_k \right| + \sup_{b,c \in B(r)} \left| [D_x T_\nu(\bar{x} + b)c]_k \right| - \frac{r}{\omega_k^s} \leq p_k(r)$$

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Lemma: If there exists $r > 0$ such that $p_k(r) < 0$ for all k , then there is a unique $\hat{x} \in B_{\bar{x}}(r)$ s.t. $f(\hat{x}, \nu) = 0$.

proof. Banach fixed point theorem.

Analytic estimates to construct the polynomials

Suppose there exist A_1, A_2, \dots, A_n such that for every $j \in \{1, \dots, n\}$ and every $\mathbf{k} \in \mathbb{Z}^d$, we have that

$$|c_{\mathbf{k}}^{(j)}| \leq \frac{A_j}{\omega_{\mathbf{k}}^{s_j}},$$

$$\omega_{\mathbf{k}}^{\mathbf{s}} = |k_1|^{s_1} \cdots |k_d|^{s_d}$$

Then, for any $\mathbf{k} \in \mathbb{Z}^d$, we get that

$$\left| \left(c^{(1)} * \cdots * c^{(n)} \right)_{\mathbf{k}} \right| \leq \left(\prod_{j=1}^n A_j \right) \frac{\alpha_{\mathbf{k}}^{(n)}}{\omega_{\mathbf{k}}^{\mathbf{s}}}.$$

Proof.

$$\begin{aligned} \left| \left(c^{(1)} * \cdots * c^{(n)} \right)_{\mathbf{k}} \right| &= \left| \sum_{\substack{\mathbf{k}^1 + \cdots + \mathbf{k}^n = \mathbf{k} \\ \mathbf{k}^1, \dots, \mathbf{k}^n \in \mathbb{Z}^d}} c_{\mathbf{k}^1}^{(1)} \cdots c_{\mathbf{k}^n}^{(n)} \right| \leq \sum_{\substack{\mathbf{k}^1 + \cdots + \mathbf{k}^n = \mathbf{k} \\ \mathbf{k}^1, \dots, \mathbf{k}^n \in \mathbb{Z}^d}} \frac{A_1}{\omega_{\mathbf{k}^1}^{s_1}} \cdots \frac{A_n}{\omega_{\mathbf{k}^n}^{s_n}} \\ &= \left(\prod_{j=1}^n A_j \right) \left(\sum_{\substack{\mathbf{k}^1 + \cdots + \mathbf{k}^n = \mathbf{k} \\ \mathbf{k}^1, \dots, \mathbf{k}^n \in \mathbb{Z}^d}} \frac{1}{\omega_{\mathbf{k}^1}^{s_1} \cdots \omega_{\mathbf{k}^n}^{s_n}} \right) \\ &= \left(\prod_{j=1}^n A_j \right) \left(\sum_{\substack{\mathbf{k}^1 + \cdots + \mathbf{k}^n = \mathbf{k} \\ \mathbf{k}^1, \dots, \mathbf{k}^n \in \mathbb{Z}^d}} \prod_{j=1}^d \frac{1}{\omega_{k_j^1}^{s_j} \cdots \omega_{k_j^n}^{s_j}} \right) \\ &= \left(\prod_{j=1}^n A_j \right) \left(\prod_{j=1}^d \sum_{\substack{k_j^1 + \cdots + k_j^n = k_j \\ k_j^1, \dots, k_j^n \in \mathbb{Z}}} \frac{1}{\omega_{k_j^1}^{s_j} \cdots \omega_{k_j^n}^{s_j}} \right) \\ &\leq \left(\prod_{j=1}^n A_j \right) \prod_{j=1}^d \frac{\alpha_{k_j}^{(n)}}{\omega_{k_j}^{s_j}} = \left(\prod_{j=1}^n A_j \right) \frac{\alpha_{\mathbf{k}}^{(n)}}{\omega_{\mathbf{k}}^{\mathbf{s}}}. \end{aligned}$$

Radii polynomials

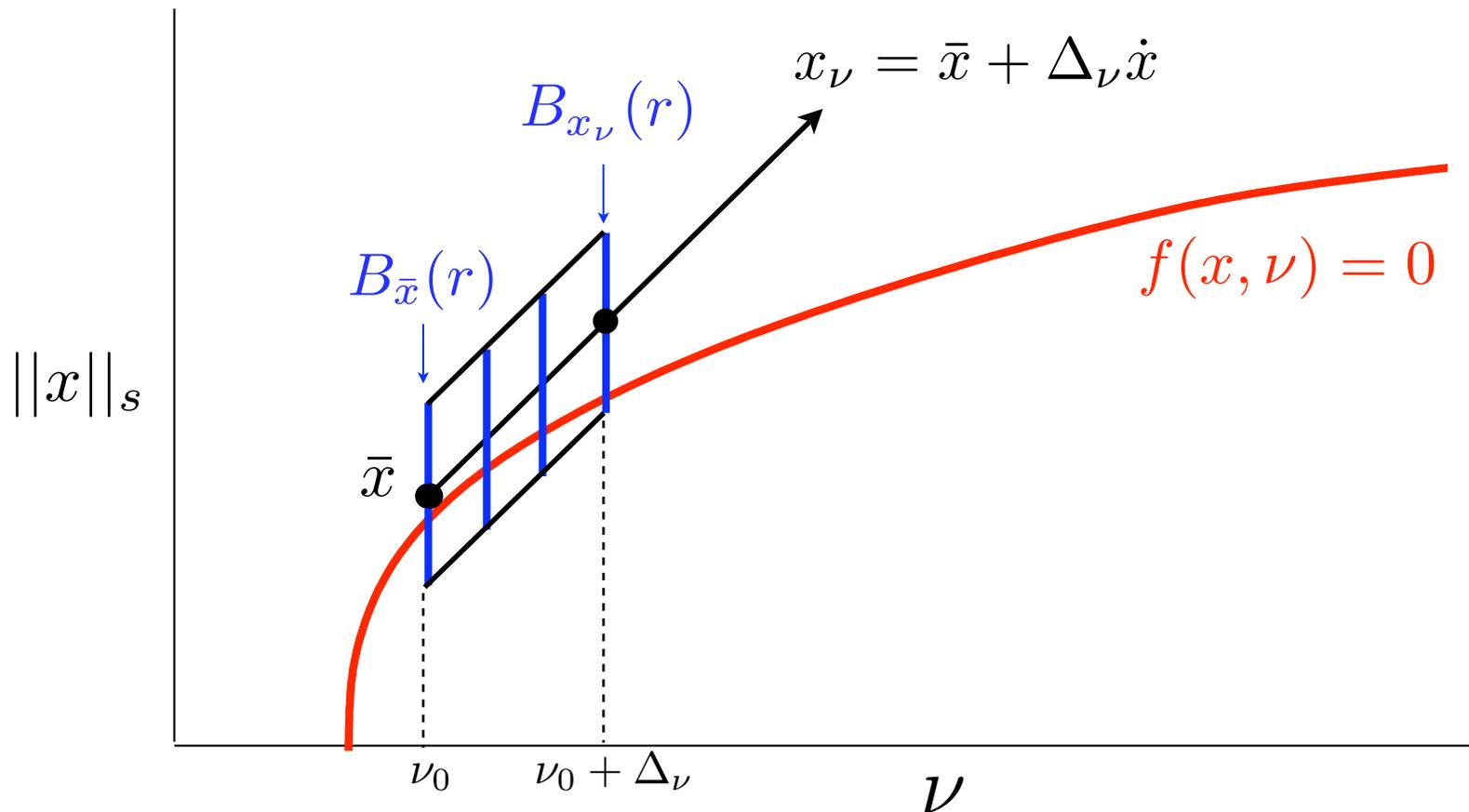
$$\{p_k(r, \Delta_\nu)\}$$



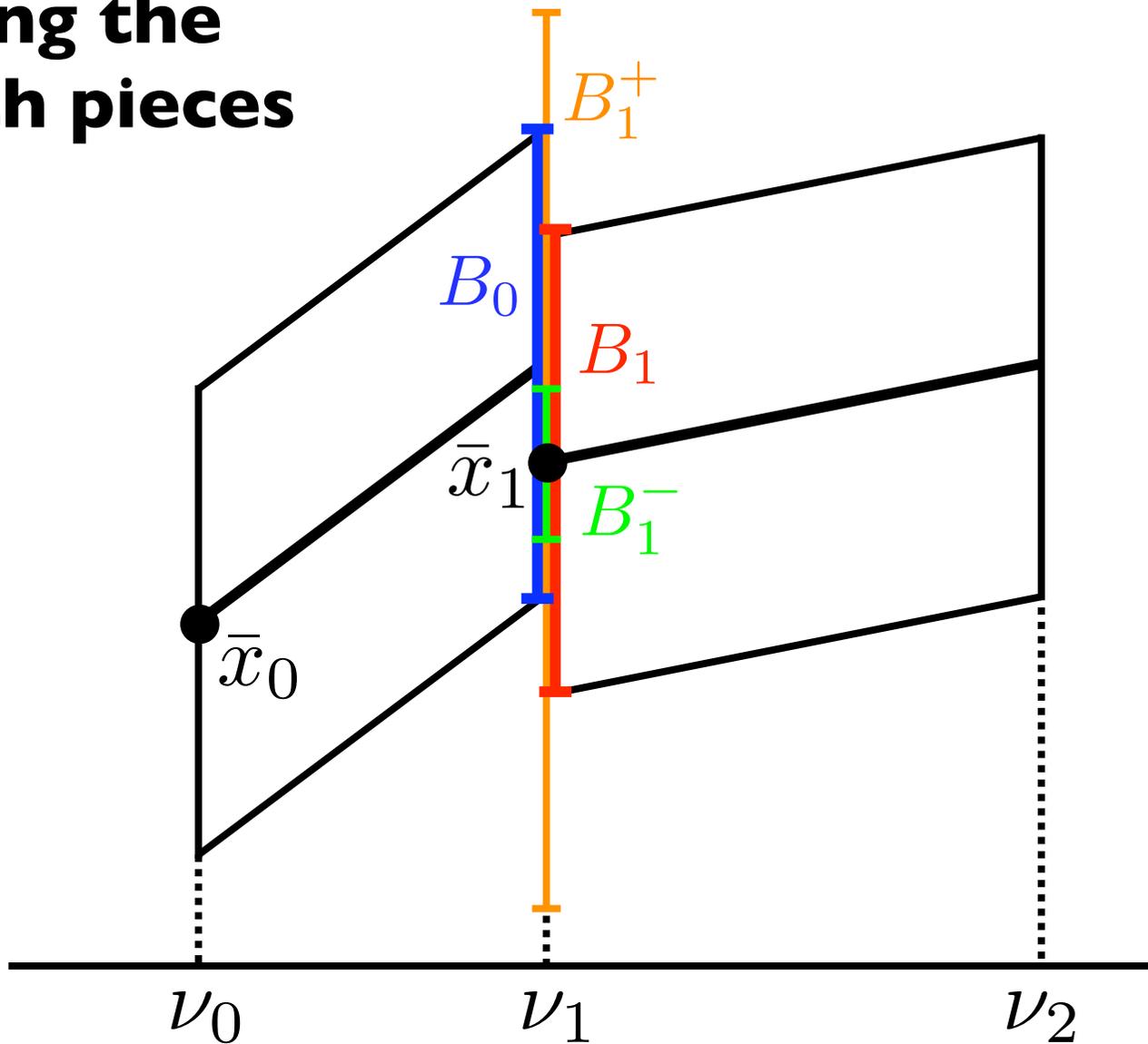
Verifying the uniform contraction principle.

$\exists r > 0$ s.t. $p_k(r, \Delta_\nu) < 0, \forall k \implies T$: uniform contraction on $[\nu_0, \nu_0 + \Delta_\nu]$

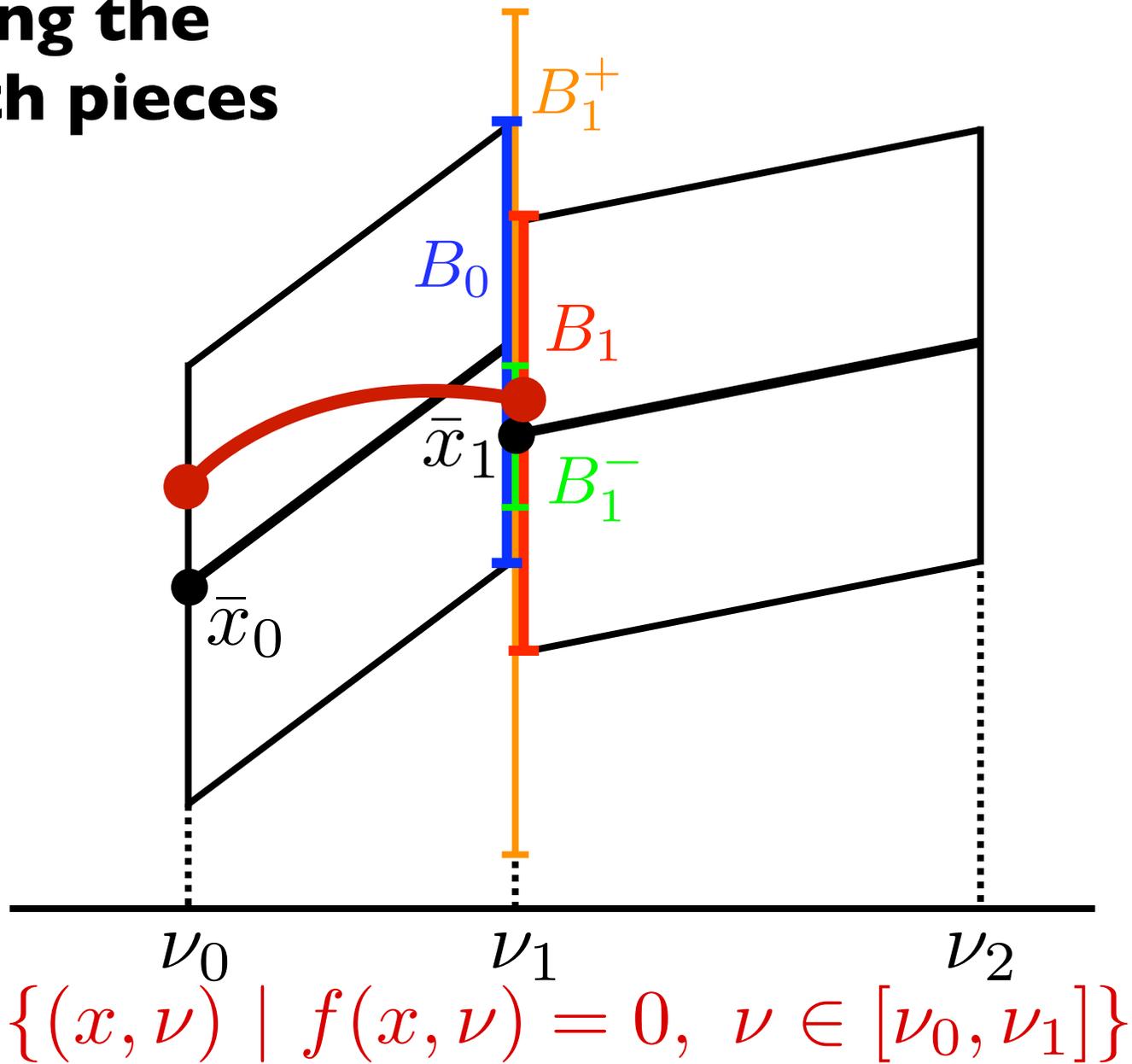
The rigorous computational method



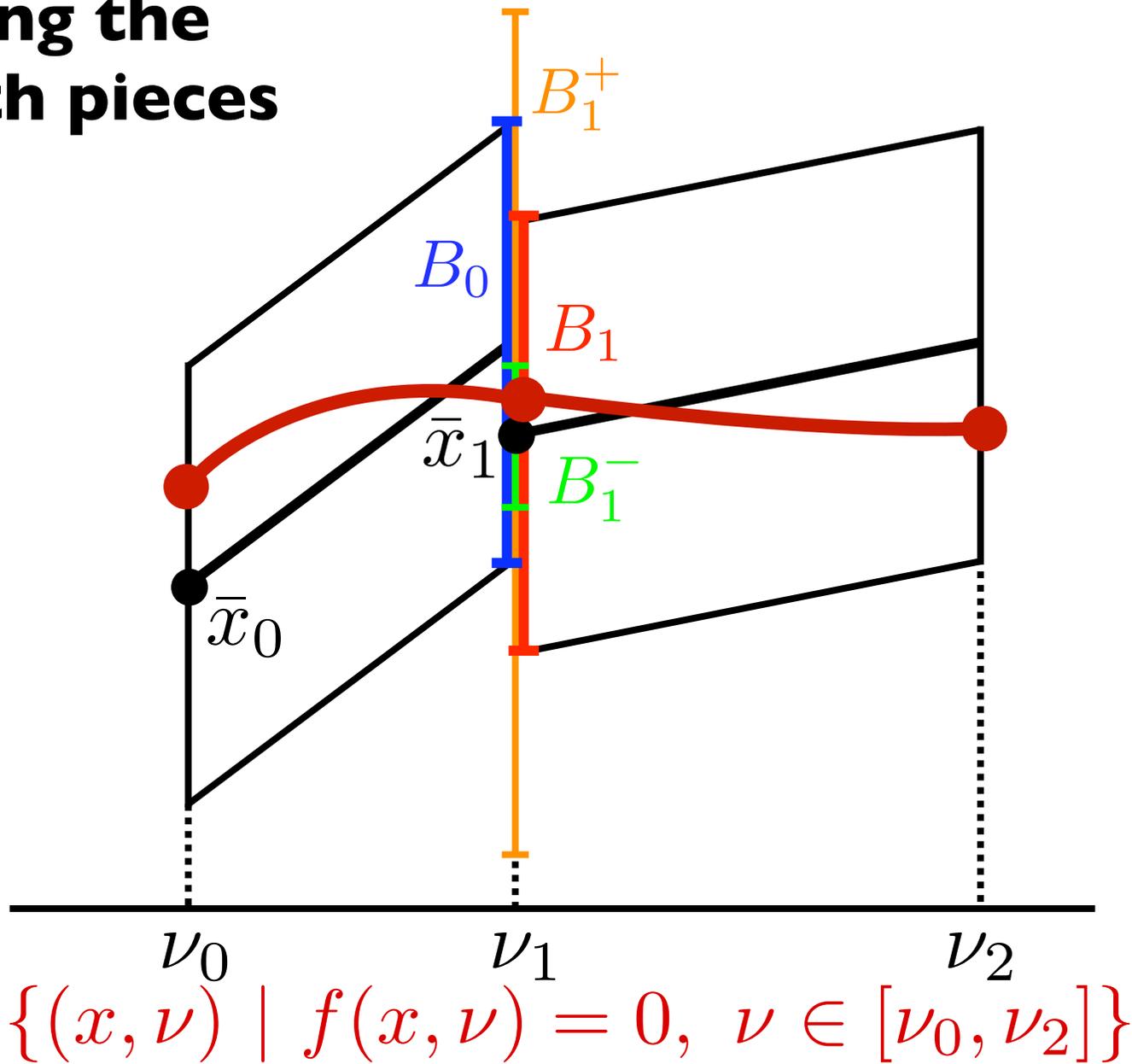
Gluing the smooth pieces



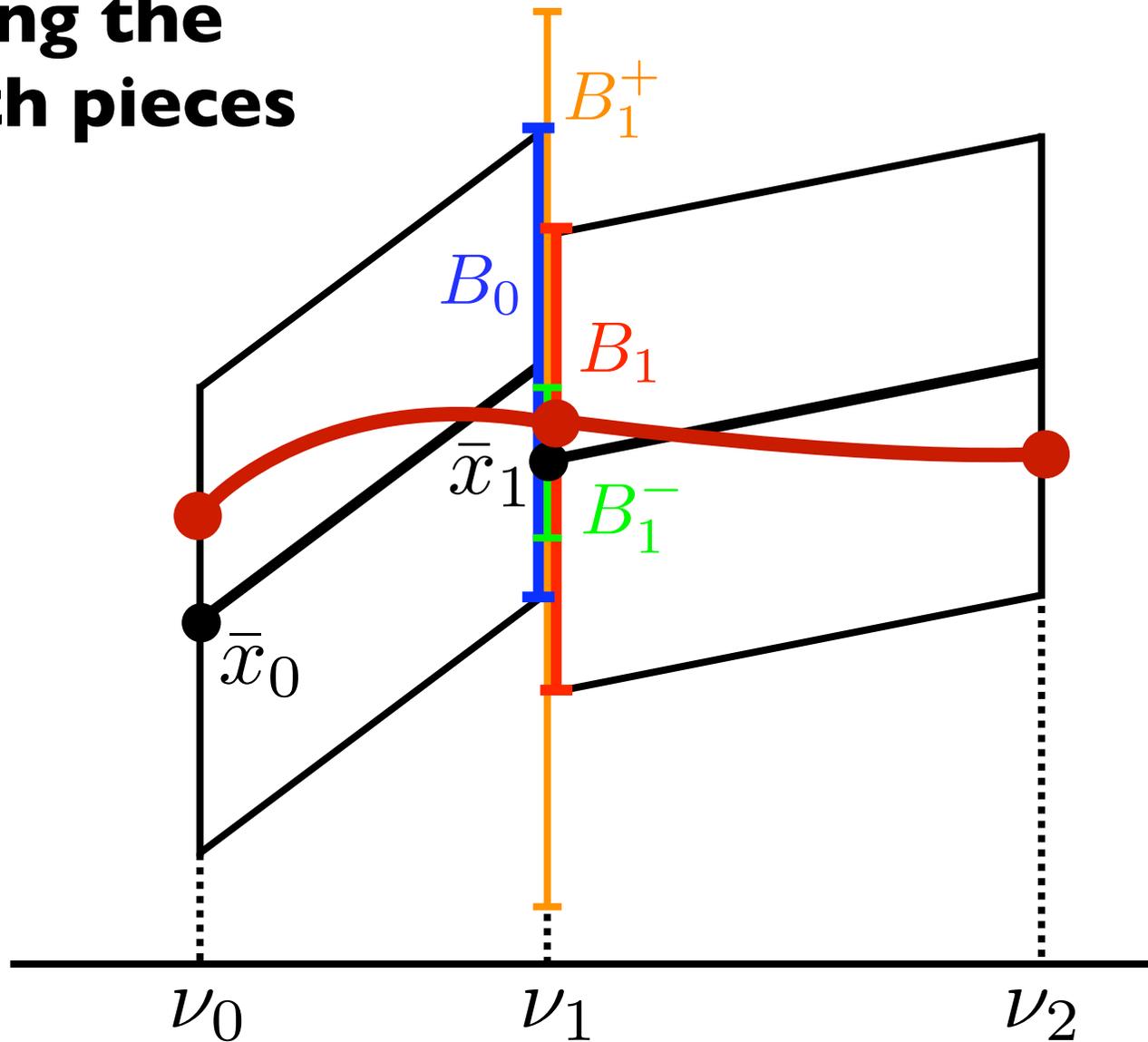
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$$\{(x, \nu) \mid f(x, \nu) = 0, \nu \in [\nu_0, \nu_2]\}$$

- **Global smooth curves of solutions.**
- **Local uniqueness by the Banach fixed point theorem.**
- **Proof of non existence of secondary bifurcations along the curves.**

Applications

- **Initial value problems of ODEs (Chebyshev in time)**
- **Boundary value problems of ODEs (Chebyshev in time)**
- **Periodic solutions of ODEs (Fourier in time)**
- **Connecting orbits of ODEs (Chebyshev in time + parameterization of invariant manifolds using power series)**
- **Equilibria of PDEs (Fourier in space)**
- **Periodic solutions of delay differential equations (Fourier in time)**
- **Minimizers of action functionals (Chebyshev in time)**
- **Periodic solutions of PDEs (Fourier in space and in time)**

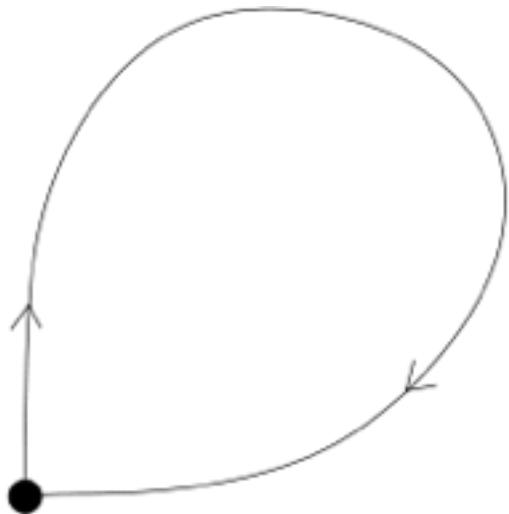
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I. Homoclinic and heteroclinic orbits of ODEs (traveling waves)

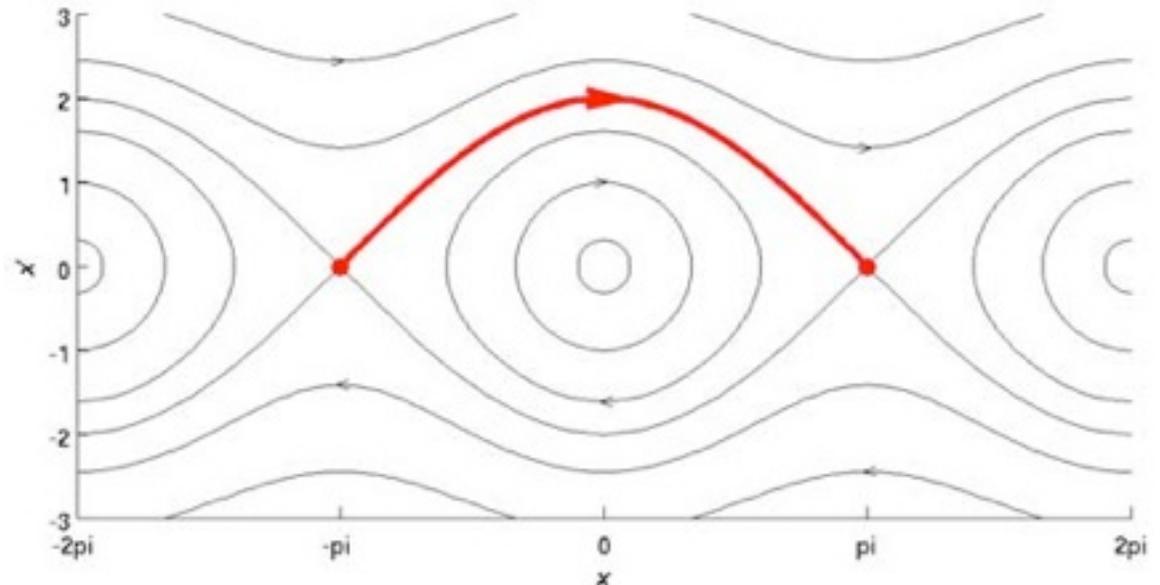
ODEs $\frac{dx}{dt} = f(x)$

$$\lim_{t \rightarrow \pm\infty} x(t) = x^{\pm} \in \mathbb{R}^n$$



$$x^{+} = x^{-}$$

homoclinic orbit

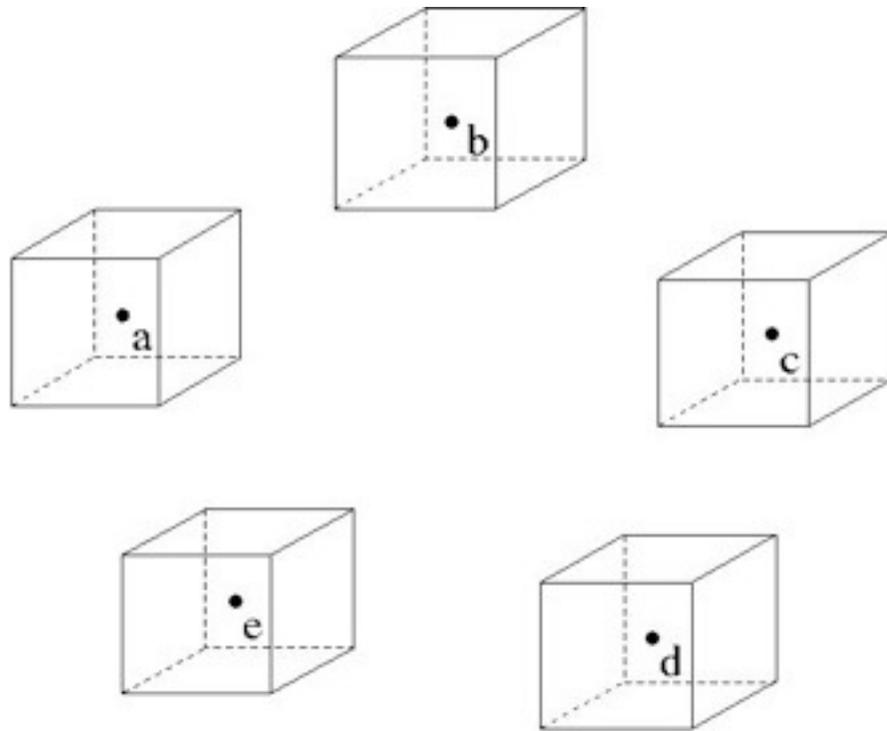


$$x^{+} \neq x^{-}$$

heteroclinic orbit

Rigorous Computations

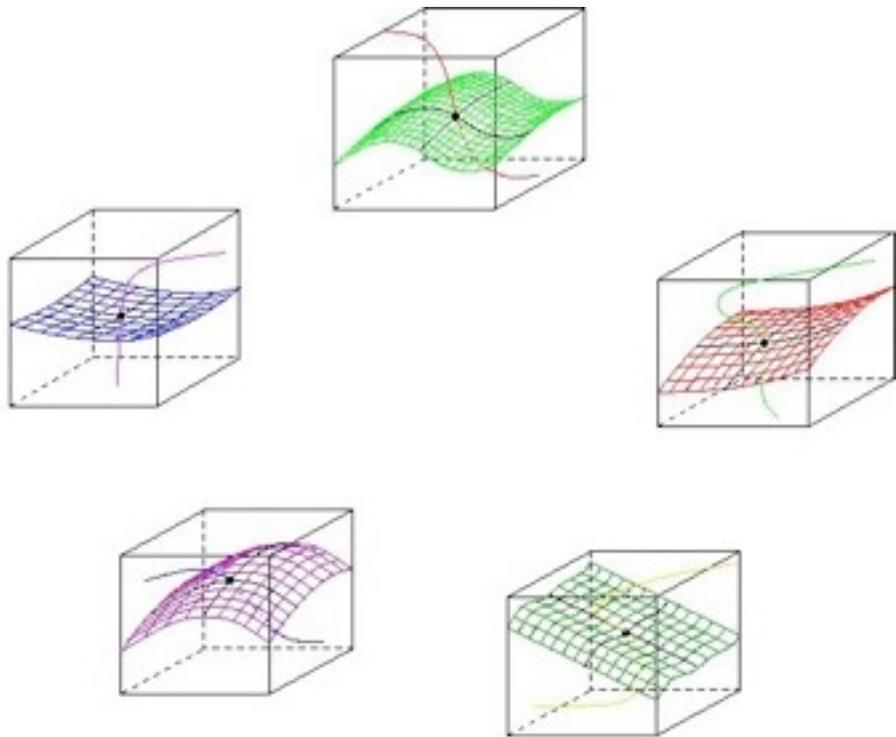
Connecting Orbits



Compute a set of equilibria.

Rigorous Computations

Connecting Orbits



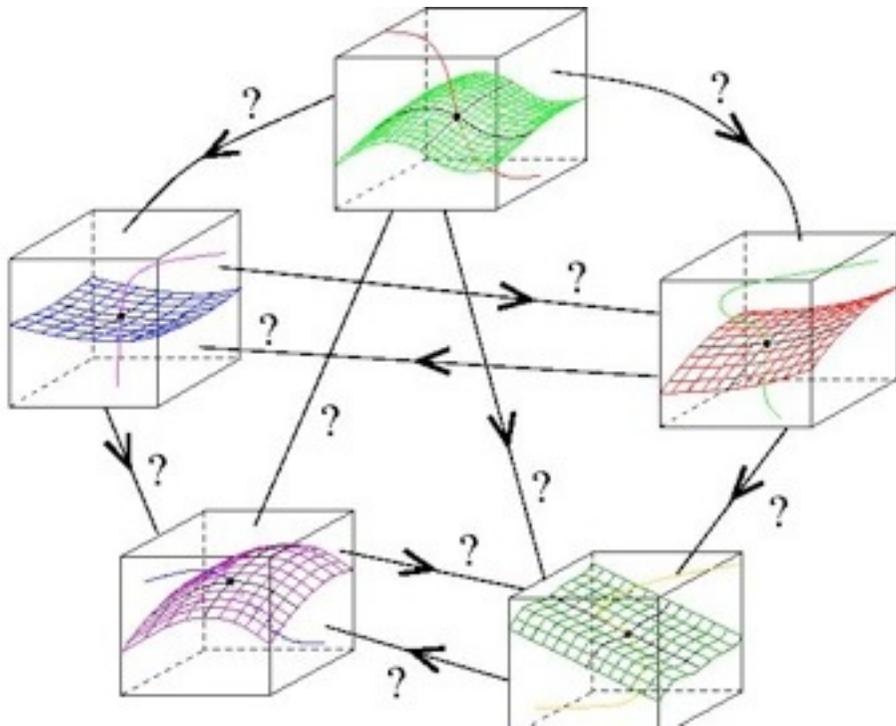
Compute a set of equilibria.

Local representation of the invariant manifolds.

Parameterization method

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Connecting orbits between the equilibria?

Boundary value problem

Chebyshev series

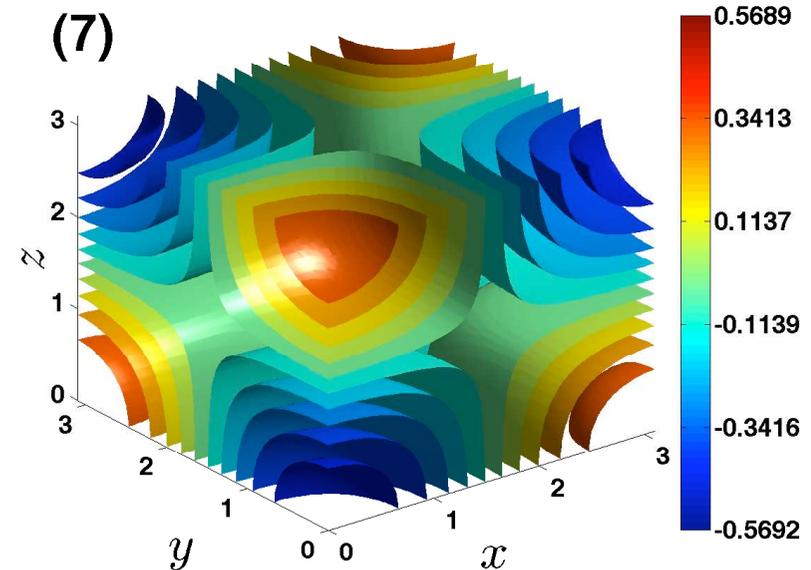
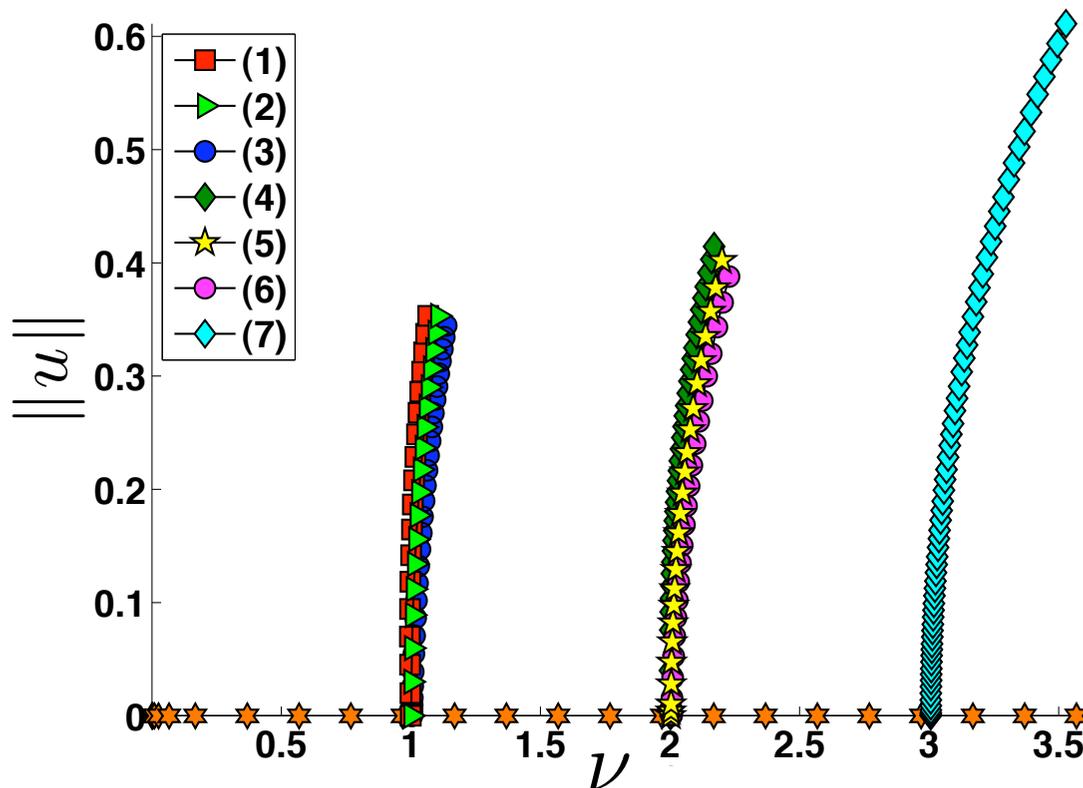
Radii polynomials

2. Equilibria of PDEs

Cahn-Hilliard 3D

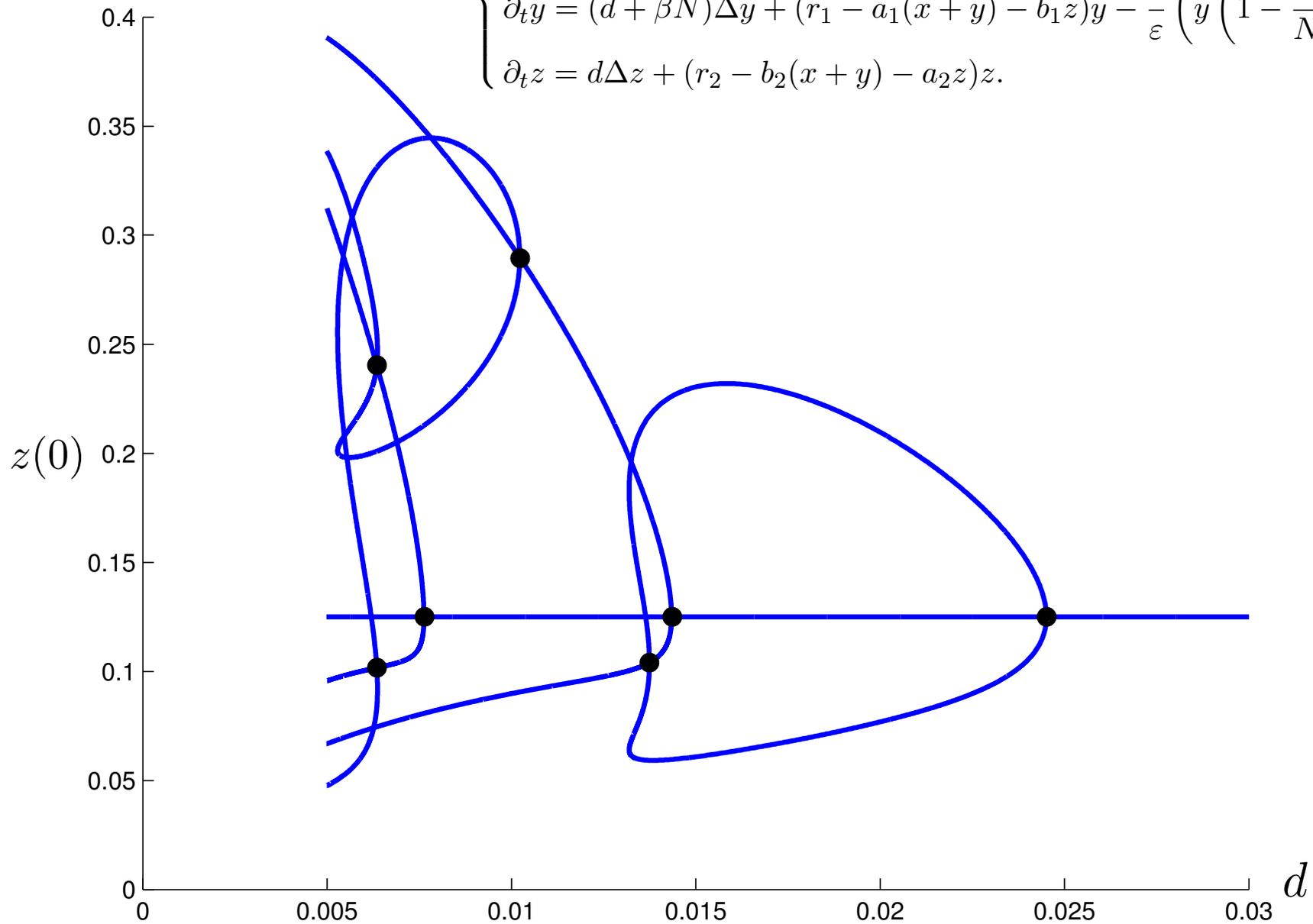
$$\begin{cases} u_t = -\Delta\left(\frac{1}{\nu}\Delta u + u - u^3\right), & \text{in } \Omega \\ \frac{\partial u}{\partial n} = \frac{\partial \Delta u}{\partial n} = 0, & \text{on } \partial\Omega \end{cases}$$

$$\Omega = [0, \pi] \times [0, \frac{\pi}{1.001}] \times [0, \frac{\pi}{1.002}]$$



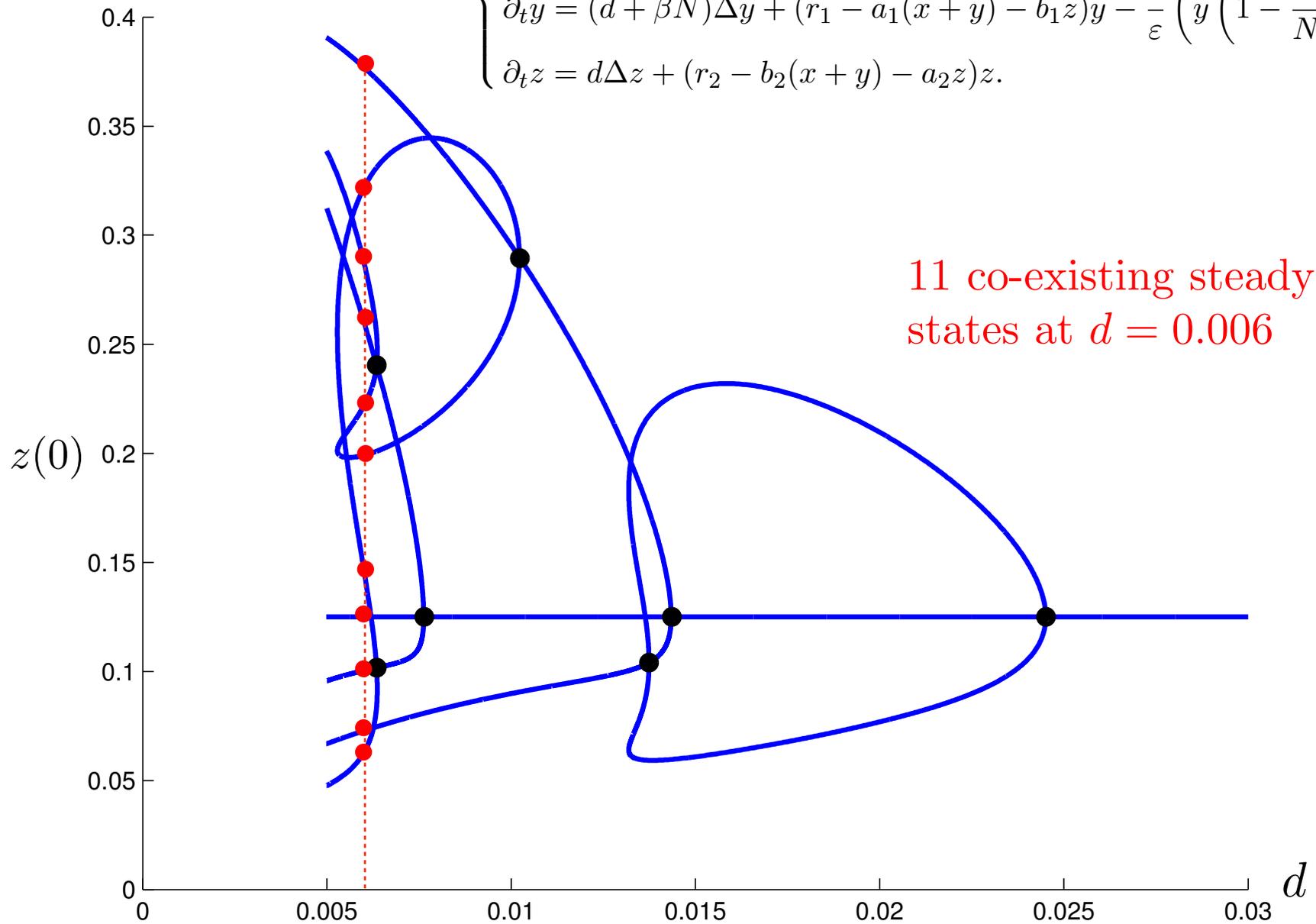
Systems of reaction-diffusion PDEs

$$\begin{cases} \partial_t x = d\Delta x + (r_1 - a_1(x + y) - b_1 z)x + \frac{1}{\varepsilon} \left(y \left(1 - \frac{z}{N} \right) - x \frac{z}{N} \right), \\ \partial_t y = (d + \beta N)\Delta y + (r_1 - a_1(x + y) - b_1 z)y - \frac{1}{\varepsilon} \left(y \left(1 - \frac{z}{N} \right) - x \frac{z}{N} \right), \\ \partial_t z = d\Delta z + (r_2 - b_2(x + y) - a_2 z)z. \end{cases}$$



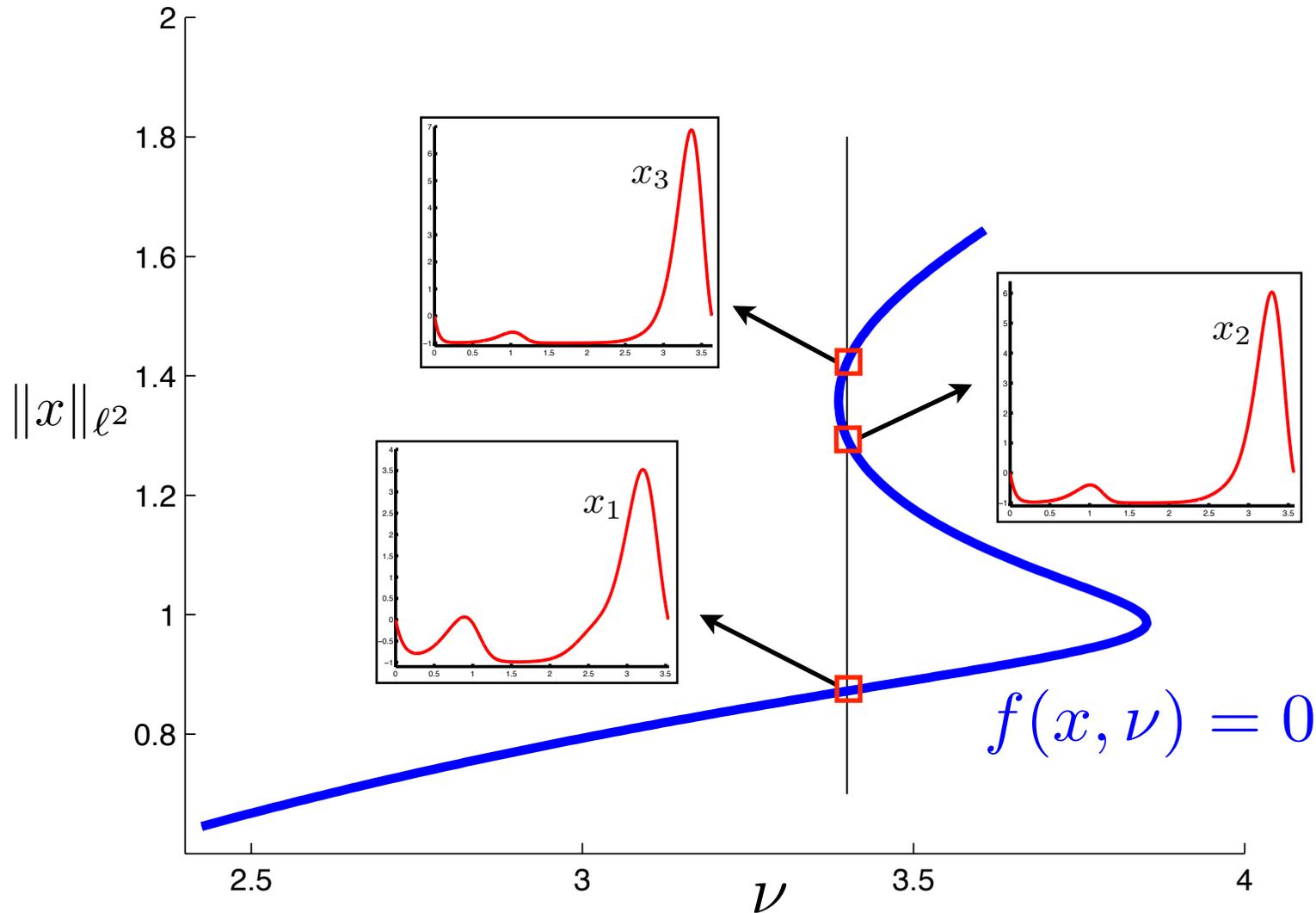
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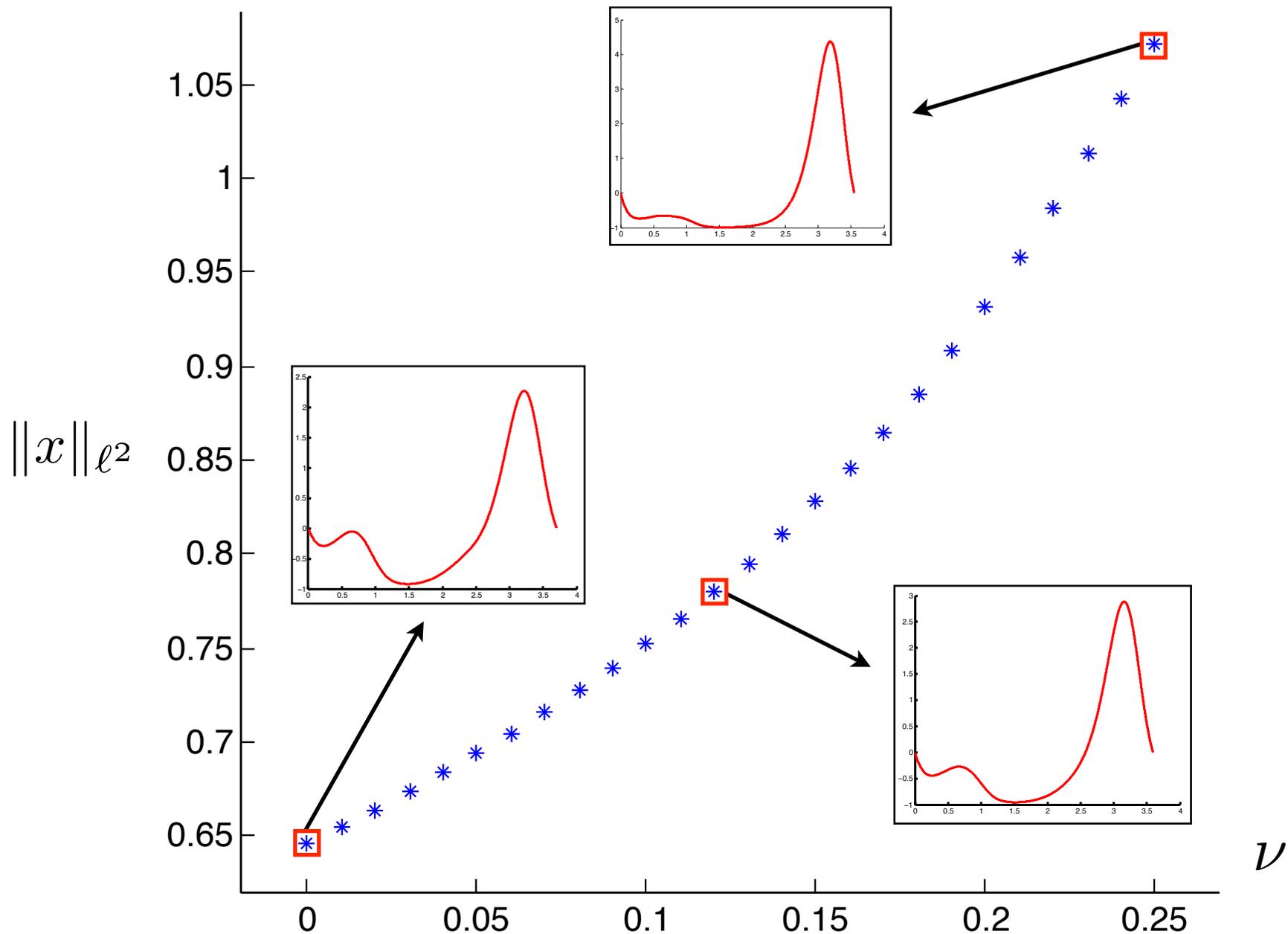


3. Periodic solutions of delay equations

$$y'(t) = \mathcal{F}(y(t), y(t - \tau_1), \dots, y(t - \tau_d)),$$



$$y'(t) = -\nu [y(t - \tau_1) + y(t - \tau_2)] [1 + y(t)],$$



$$y'(t) = - [2.425y(t - \tau_1) + 2.425y(t - \tau_2) + \nu y(t - \tau_3)] [1 + y(t)],$$

4. Minimizers of action functionals

Ginzburg–Landau energy: a model of superconductivity

$$G = G(\phi, a) = \frac{1}{2d} \int_{-d}^d \left(\phi^2(\phi^2 - 2) + \frac{2(\phi')^2}{\kappa^2} + 2\phi^2 a^2 + 2(a' - h_e)^2 \right) dt.$$

$\phi > 0$: measures the density of superconducting electrons

a : magnetic field potential

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Parameters

d : size of the superconducting material

h_e : external magnetic field

κ : Ginzburg-Landau parameter.

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$\kappa = 0.3, d = 4$

Parameters

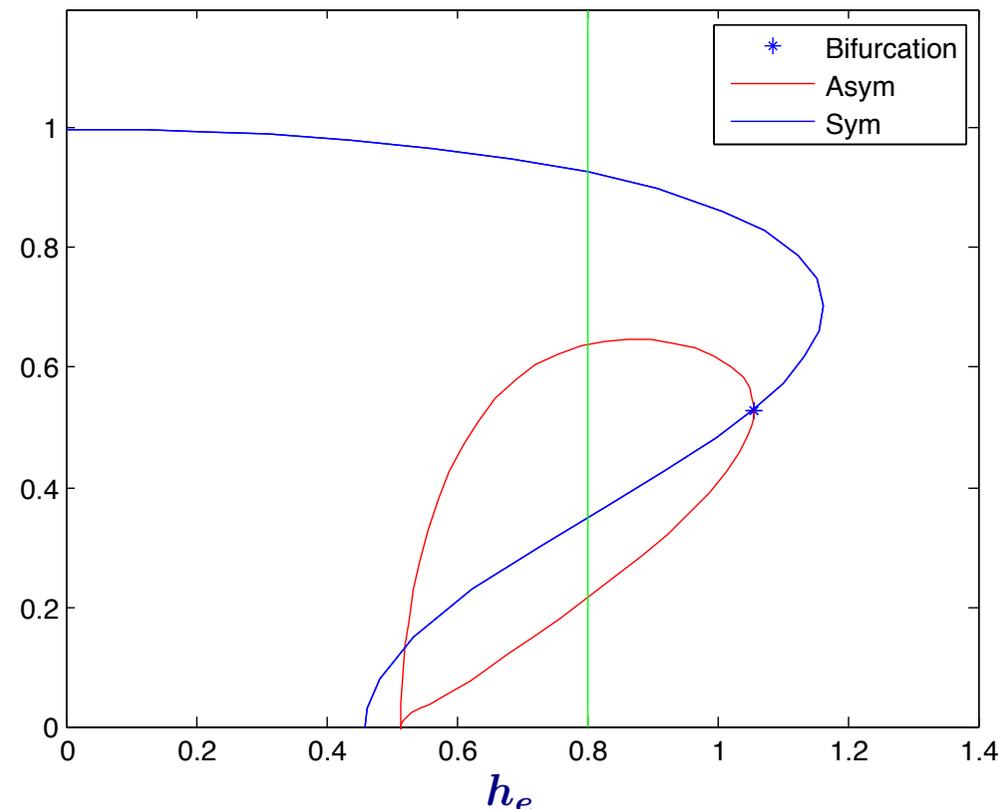
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$\phi(d)$

**Co-existence of
nontrivial solutions**



5. Periodic orbits of PDEs

Kuramoto-Sivashinski equation

$$(KS) \begin{cases} u_t = -\nu u_{yyyyy} - u_{yy} + 2uu_y \\ u(t, y) = u(t, y + 2\pi), \quad u(t, -y) = -u(t, y) \end{cases}$$

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Goal: propose an method (based on spectral methods and fixed point theory) to rigorously compute time periodic solutions of PDEs.

Letting $L = \frac{2\pi}{p}$, the time-periodic solutions of period p of (KS) can be expanded using the Fourier expansion

$$u(t, y) = \sum_{\mathbf{k} \in \mathbb{Z}^2} c_{\mathbf{k}} \psi_{\mathbf{k}}, \quad \text{where for } \mathbf{k} = (k_1, k_2) \in \mathbb{Z}^2, \quad \psi_{\mathbf{k}} = e^{iLk_1 t} e^{ik_2 y}.$$

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$$x_{\mathbf{k}} = \begin{cases} L, & \mathbf{k} = (0, 0) \\ b_{\mathbf{k}}, & \mathbf{k} = (0, k_2), \quad k_2 \neq 0 \\ \begin{pmatrix} a_{\mathbf{k}} \\ b_{\mathbf{k}} \end{pmatrix}, & \mathbf{k} = (k_1, k_2), \quad k_1 \neq 0 \text{ and } k_2 \neq 0. \end{cases}$$

$$a_{\mathbf{k}} \stackrel{\text{def}}{=} \text{Re}(c_{\mathbf{k}}) \quad \text{and} \quad b_{\mathbf{k}} \stackrel{\text{def}}{=} \text{Im}(c_{\mathbf{k}}).$$

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Plugging the space-time Fourier expansion into (KS) results in solving, for all $\mathbf{k} \in \mathbb{Z}^2$

$$h_{\mathbf{k}} \stackrel{\text{def}}{=} \mu_{\mathbf{k}} c_{\mathbf{k}} - 2 \sum_{\mathbf{k}^1 + \mathbf{k}^2 = \mathbf{k}} ik_2^1 c_{\mathbf{k}^1} c_{\mathbf{k}^2} = \mu_{\mathbf{k}} c_{\mathbf{k}} - k_2 i \sum_{\mathbf{k}^1 + \mathbf{k}^2 = \mathbf{k}} c_{\mathbf{k}^1} c_{\mathbf{k}^2} = 0,$$

where $\mu_{\mathbf{k}} = \mu_{k_1, k_2} \stackrel{\text{def}}{=} ik_1 L + \nu k_2^4 - k_2^2$.

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$$f_{\mathbf{k}} \stackrel{\text{def}}{=} \text{Re}(h_{\mathbf{k}}) = (\nu k_2^4 - k_2^2) a_{\mathbf{k}} - (k_1 L) b_{\mathbf{k}} + 2k_2 \sum_{\mathbf{k}^1 + \mathbf{k}^2 = \mathbf{k}} a_{\mathbf{k}^1} b_{\mathbf{k}^2},$$

$$g_{\mathbf{k}} \stackrel{\text{def}}{=} \text{Im}(h_{\mathbf{k}}) = (k_1 L) a_{\mathbf{k}} + (\nu k_2^4 - k_2^2) b_{\mathbf{k}} - k_2 \sum_{\mathbf{k}^1 + \mathbf{k}^2 = \mathbf{k}} (a_{\mathbf{k}^1} a_{\mathbf{k}^2} - b_{\mathbf{k}^1} b_{\mathbf{k}^2})$$

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Functions

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Unknowns

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$$\mathcal{I} = \{(0, 0)\} \cup \{\mathbf{k} = (0, k_2) \mid k_2 \neq 0\} \cup \{\mathbf{k} = (k_1, k_2) \mid k_1 \neq 0 \text{ and } k_2 \neq 0\},$$

one can identify $x = \{x_{\mathbf{k}}\}_{\mathbf{k} \in \mathcal{I}}$.

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Finally, let us define $\mathcal{F} = \{\mathcal{F}_{\mathbf{k}}\}_{\mathbf{k} \in \mathcal{I}}$ component-wise by

$$\mathcal{F}_{\mathbf{k}} = \begin{cases} \eta, & \mathbf{k} = (0, 0) \\ g_{\mathbf{k}}, & \mathbf{k} = (0, k_2), \quad k_2 \neq 0 \\ \begin{pmatrix} f_{\mathbf{k}} \\ g_{\mathbf{k}} \end{pmatrix}, & \mathbf{k} = (k_1, k_2), \quad k_1 \neq 0 \text{ and } k_2 \neq 0. \end{cases}$$

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Lemma. Finding time-periodic solutions $u(t, y)$ of (KS) such that $\eta = 0$ is equivalent to find x such that $\mathcal{F}(x) = 0$.

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**To solve rigorously
in a Banach space**

The Banach space

Define the one-dimensional weights ω_k^s by

$$\omega_k^s \stackrel{\text{def}}{=} \begin{cases} 1, & \text{if } k = 0 \\ |k|^s, & \text{if } k \neq 0. \end{cases}$$

Using the 1-d weights, define the 2-dimensional weights, given $\mathbf{k} = (k_1, k_2) \in \mathbb{Z}^2$,

$$\omega_{\mathbf{k}}^{\mathbf{s}} \stackrel{\text{def}}{=} \omega_{k_1}^{s_1} \omega_{k_2}^{s_2}.$$

They are used to define the norm

$$\|x\|_{\mathbf{s}} = \sup_{\mathbf{k} \in \mathcal{I}} \omega_{\mathbf{k}}^{\mathbf{s}} |x_{\mathbf{k}}|_{\infty},$$

where $|x_{\mathbf{k}}|_{\infty}$ is the sup norm of the vector $x_{\mathbf{k}}$, which is one or two dimensional, depending on \mathbf{k} . Define the Banach space

$$X^{\mathbf{s}} = \{x \mid \|x\|_{\mathbf{s}} < \infty\},$$

consisting of sequences with algebraically decaying tails according to the rate \mathbf{s} .

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**Banach algebra
under discrete
convolution**

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For sake of simplicity of the presentation, for $\mathbf{k} = (k_1, k_2)$ with $k_1 \neq 0$ or $k_2 \neq 0$, let

$$R_{\mathbf{k}}(\nu, L) \stackrel{\text{def}}{=} \begin{pmatrix} \nu k_2^4 - k_2^2 & -k_1 L \\ k_1 L & \nu k_2^4 - k_2^2 \end{pmatrix} \quad \text{and} \quad R_{0, k_2}(\nu, L) \stackrel{\text{def}}{=} \nu k_2^4 - k_2^2,$$

$$\mathcal{N}_{\mathbf{k}}(x) \stackrel{\text{def}}{=} \sum_{\mathbf{k}^1 + \mathbf{k}^2 = \mathbf{k}} \begin{pmatrix} 2a_{\mathbf{k}^1} b_{\mathbf{k}^2} \\ -a_{\mathbf{k}^1} a_{\mathbf{k}^2} + b_{\mathbf{k}^1} b_{\mathbf{k}^2} \end{pmatrix}$$

so that one has that

$$\mathcal{F}_{\mathbf{k}}(x, \nu) = R_{\mathbf{k}}(\nu, L)x_{\mathbf{k}} + k_2 \mathcal{N}_{\mathbf{k}}(x).$$

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$$\mathcal{N}_{\mathbf{k}}(x) \stackrel{\text{def}}{=} \sum_{\mathbf{k}^1 + \mathbf{k}^2 = \mathbf{k}} \begin{pmatrix} 2a_{\mathbf{k}^1} b_{\mathbf{k}^2} \\ -a_{\mathbf{k}^1} a_{\mathbf{k}^2} + b_{\mathbf{k}^1} b_{\mathbf{k}^2} \end{pmatrix}$$

so that one has that

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Lemma. (Bootstrap) Consider a fixed decay rate $\mathbf{s} > (1, 1)$ and assume the existence of $\mathbf{M} > (0, 0)$ such that $R_{\mathbf{k}}(\nu, L)$ is invertible for all $|\mathbf{k}| > \mathbf{M}$. If there exists $x \in X^{\mathbf{s}}$ such that $\mathcal{F}(x) = 0$, then $x \in X^{\mathbf{s}_0}$, for all $\mathbf{s}_0 > (1, 1)$.

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Hence, we focus our attention on looking for zeros of \mathcal{F} within a Banach space with a fixed decay rate $\mathbf{s} > (1, 1)$.

Given $\mathbf{m} = (m_1, m_2)$, define $\mathbf{F}_{\mathbf{m}} = F_{m_1} \times F_{m_2}$, where $F_{m_j} \stackrel{\text{def}}{=} \{k_j \in \mathbb{Z} \mid |k_j| < m_j\}$. Consider a *Galerkin projection* of \mathcal{F} of dimension $n = n(\mathbf{m}) \stackrel{\text{def}}{=} 2m_1m_2 - 2m_1 - m_2 + 2$ given by $\mathcal{F}^{(\mathbf{m})} \stackrel{\text{def}}{=} \{\mathcal{F}_{\mathbf{k}}^{(\mathbf{m})}\}_{\mathbf{k} \in \mathbf{F}_{\mathbf{m}}}$, where $\mathcal{F}^{(\mathbf{m})}: \mathbb{R}^n \rightarrow \mathbb{R}^n$, is given component-wise by

$$\mathcal{F}_{\mathbf{k}}^{(\mathbf{m})}(x_{\mathbf{F}_{\mathbf{m}}}) \stackrel{\text{def}}{=} \mathcal{F}_{\mathbf{k}}(x_{\mathbf{F}_{\mathbf{m}}}, 0_{I_{\mathbf{m}}}), \quad \mathbf{k} \in \mathbf{F}_{\mathbf{m}}.$$

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Consider $\hat{x}_{\mathbf{F}_{\mathbf{m}}}$ such that $\mathcal{F}^{(\mathbf{m})}(\hat{x}_{\mathbf{F}_{\mathbf{m}}}) \approx 0$. Let $\hat{x} \stackrel{\text{def}}{=} (\hat{x}_{\mathbf{F}_{\mathbf{m}}}, 0_{\mathbf{I}_{\mathbf{m}}}) \in X^{\mathbf{s}}$. Assume that the Jacobian matrix $D\mathcal{F}^{(\mathbf{m})}(\hat{x}_{\mathbf{F}_{\mathbf{m}}})$ is non-singular and let $A_{\mathbf{m}}$ an approximation for its inverse.

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Define the action of the linear operator A on $x = \{x_{\mathbf{k}}\}_{\mathbf{k} \in \mathcal{I}}$ component-wise by

$$\left[A(x) \right]_{\mathbf{k}} \stackrel{\text{def}}{=} \begin{cases} \left[A_{\mathbf{m}}(x_{\mathbf{F}_{\mathbf{m}}}) \right]_{\mathbf{k}}, & \text{if } \mathbf{k} \in \mathbf{F}_{\mathbf{m}} \\ R_{\mathbf{k}}(\nu, \hat{L})^{-1} x_{\mathbf{k}}, & \text{if } \mathbf{k} \notin \mathbf{F}_{\mathbf{m}}. \end{cases}$$

$$T(x) \stackrel{\text{def}}{=} x - A\mathcal{F}(x).$$

Given $\mathbf{m} = (m_1, m_2)$, define $\mathbf{F}_{\mathbf{m}} = F_{m_1} \times F_{m_2}$, where $F_{m_j} \stackrel{\text{def}}{=} \{k_j \in \mathbb{Z} \mid |k_j| < m_j\}$. Consider a *Galerkin projection* of \mathcal{F} of dimension $n = n(\mathbf{m}) \stackrel{\text{def}}{=} 2m_1m_2 - 2m_1 - m_2 + 2$ given by $\mathcal{F}^{(\mathbf{m})} \stackrel{\text{def}}{=} \{\mathcal{F}_{\mathbf{k}}^{(\mathbf{m})}\}_{\mathbf{k} \in \mathbf{F}_{\mathbf{m}}}$, where $\mathcal{F}^{(\mathbf{m})}: \mathbb{R}^n \rightarrow \mathbb{R}^n$, is given component-wise by

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(Newton-like operator)

Given $\mathbf{m} = (m_1, m_2)$, define $\mathbf{F}_{\mathbf{m}} = F_{m_1} \times F_{m_2}$, where $F_{m_j} \stackrel{\text{def}}{=} \{k_j \in \mathbb{Z} \mid |k_j| < m_j\}$. Consider a *Galerkin projection* of \mathcal{F} of dimension $n = n(\mathbf{m}) \stackrel{\text{def}}{=} 2m_1m_2 - 2m_1 - m_2 + 2$ given by $\mathcal{F}^{(\mathbf{m})} \stackrel{\text{def}}{=} \{\mathcal{F}_{\mathbf{k}}^{(\mathbf{m})}\}_{\mathbf{k} \in \mathbf{F}_{\mathbf{m}}}$, where $\mathcal{F}^{(\mathbf{m})}: \mathbb{R}^n \rightarrow \mathbb{R}^n$, is given component-wise by

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Lemma. Consider a Galerkin projection dimension $\mathbf{m} = (m_1, m_2)$ and let $\mathbf{s} = (s_1, s_2) > (1, 1)$ a decay rate. The solutions of $\mathcal{F} = 0$ are in one to one correspondence with the fixed points of T . Also, one has that $T: X^{\mathbf{s}} \rightarrow X^{\mathbf{s}}$.

The rigorous continuation method is based on the notion of the radii polynomials, which provide a numerically efficient way to verify that the operator T is a contraction on a small closed ball $B(\hat{x}, r)$ centered at the numerical approximation \hat{x} in X^s .

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Ingredients to construct the radii polynomials

- Convolution estimates
- Interval arithmetic
- Fast Fourier transform

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Ingredients to construct the radii polynomials

- Convolution estimates
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The closed ball of radius r in $X^{\mathbf{s}}$, centered at the origin, is given by

$$B(r) \stackrel{\text{def}}{=} \prod_{\mathbf{k} \in \mathcal{I}} \left[-\frac{r}{\omega_{\mathbf{k}}^{\mathbf{s}}}, \frac{r}{\omega_{\mathbf{k}}^{\mathbf{s}}} \right]^{d(\mathbf{k})},$$

where $d(\mathbf{k}) = 1$ if $\mathbf{k} = (0, k_2)$ and $d(\mathbf{k}) = 2$ otherwise. The closed ball of radius r centered at \hat{x} is then

$$B(\hat{x}, r) \stackrel{\text{def}}{=} \hat{x} + B(r).$$

Consider now bounds $Y_{\mathbf{k}}$ and $Z_{\mathbf{k}}$ for all $\mathbf{k} \in \mathcal{I}$, such that

$$\left| [T(\hat{x}) - \hat{x}]_{\mathbf{k}} \right| \leq Y_{\mathbf{k}},$$

and

$$\sup_{x_1, x_2 \in B(r)} \left| [DT(\hat{x} + x_1)x_2]_{\mathbf{k}} \right| \leq Z_{\mathbf{k}}(r).$$

Lemma. If there exists an $r > 0$ such that $\|Y + Z\|_{\mathbf{s}} < r$, with $Y \stackrel{\text{def}}{=} \{Y_{\mathbf{k}}\}_{\mathbf{k} \in \mathcal{I}}$ and $Z \stackrel{\text{def}}{=} \{Z_{\mathbf{k}}\}_{\mathbf{k} \in \mathcal{I}}$, then T is a contraction mapping on $B(\hat{x}, r)$ with contraction constant at most $\|Y + Z\|_{\mathbf{s}}/r < 1$. Furthermore, there is a unique $\tilde{x} \in B(\hat{x}, r)$ such that $\mathcal{F}(\tilde{x}) = 0$.

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Define the *finite radii polynomials* $\{p_{\mathbf{k}}(r)\}_{\mathbf{k} \in F_M}$ by

$$p_{\mathbf{k}}(r) \stackrel{\text{def}}{=} Y_{\mathbf{k}} + Z_{\mathbf{k}}(r) - \frac{r}{\omega_{\mathbf{k}}^{\mathbf{s}}} \mathbb{I}^{d(\mathbf{k})},$$

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for $Z_{\mathbf{k}}$ in $\mathbf{X}^{\mathbf{s}}$**

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**asymptotic bound
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Lemma. If there exists $r > 0$ such that $p_{\mathbf{k}}(r) < 0$ for all $\mathbf{k} \in F_M$ and $\tilde{p}_M(r) < 0$, then there is a unique $\tilde{x} \in B(\hat{x}, r)$ such that $\mathcal{F}(\tilde{x}) = 0$.

Results

Kuramoto-Sivashinski equation

$$(KS) \begin{cases} u_t = -\nu u_{yyyyy} - u_{yy} + 2uu_y \\ u(t, y) = u(t, y + 2\pi), \quad u(t, -y) = -u(t, y) \end{cases}$$

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of time Fourier modes

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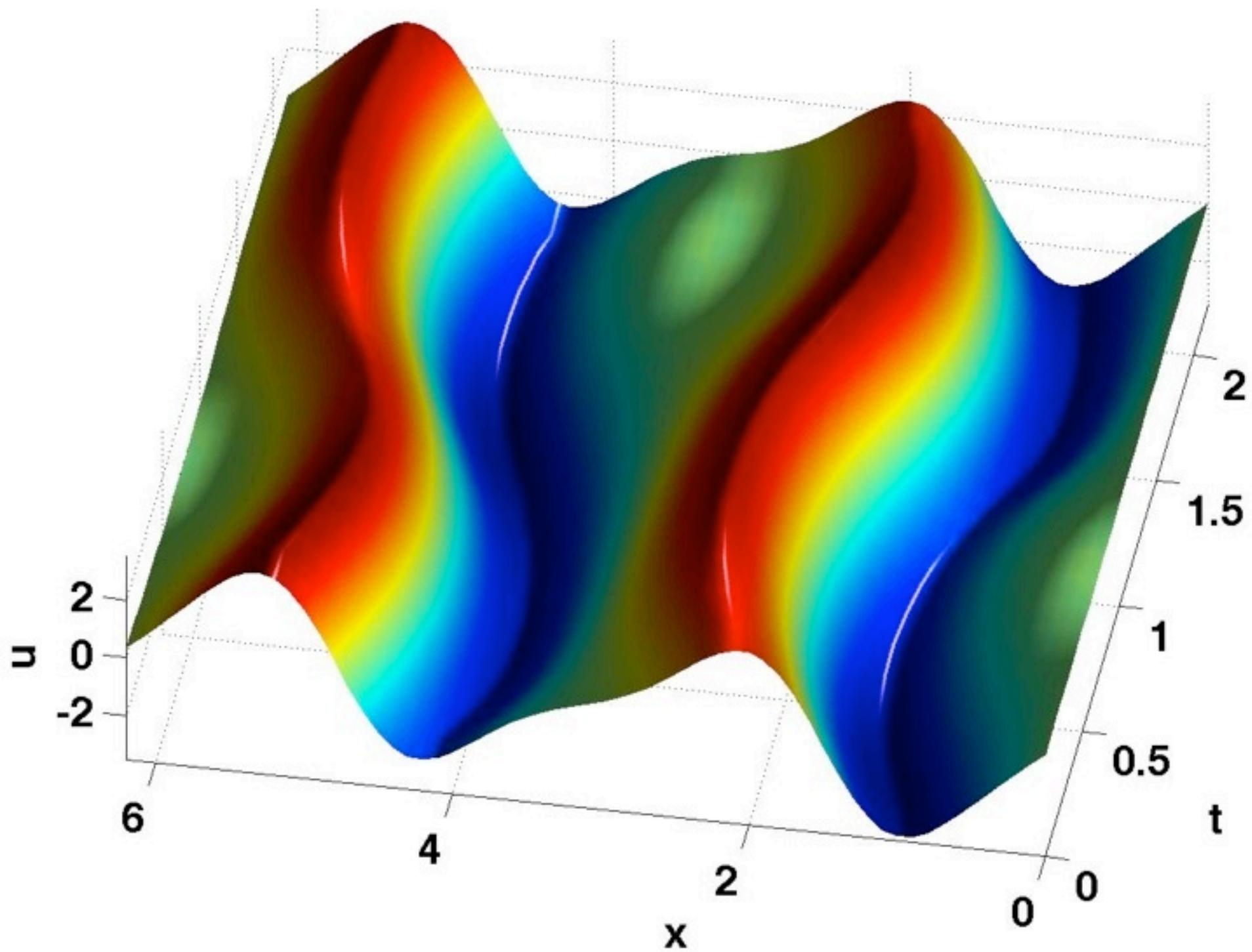
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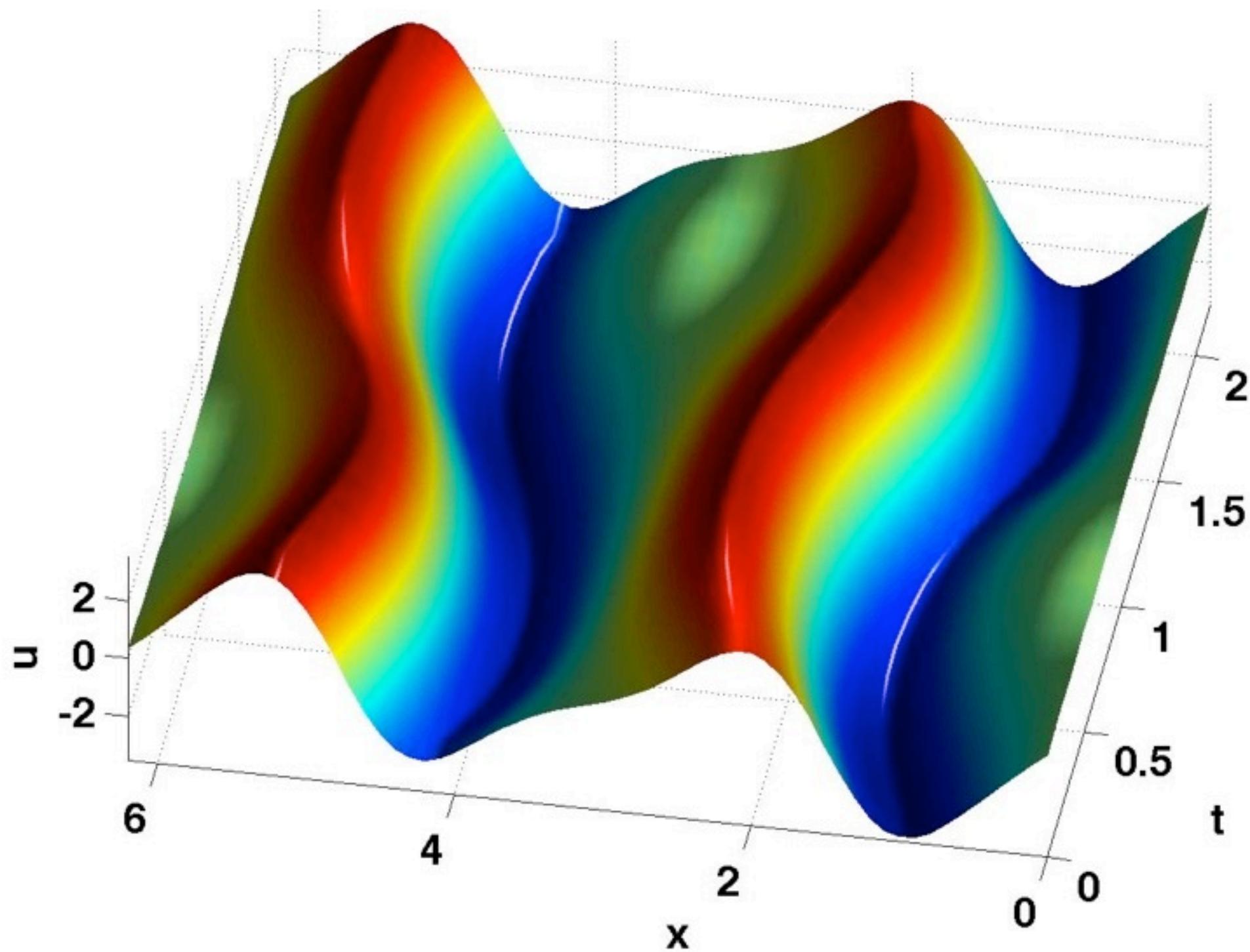
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$$\tilde{x} \in B(\hat{x}, r) = \hat{x} + \prod_{\mathbf{k} \in \mathcal{I}} \left[-\frac{3 \times 10^{-4}}{k_1^{3/2} k_2^{3/2}}, \frac{3 \times 10^{-4}}{k_1^{3/2} k_2^{3/2}} \right]^{d(\mathbf{k})} \subset X^{(\frac{3}{2}, \frac{3}{2})}$$

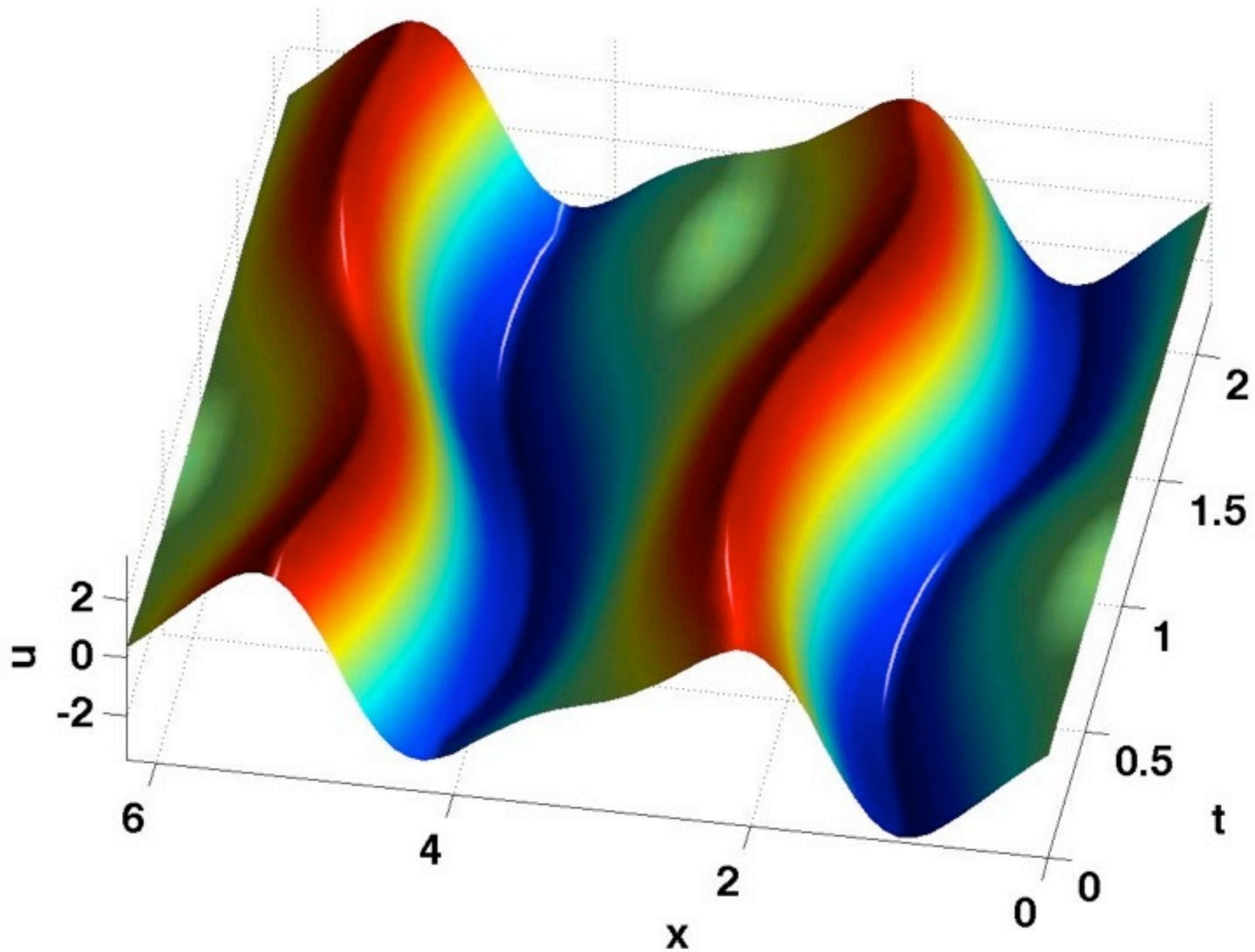
$$\nu = 0.127$$



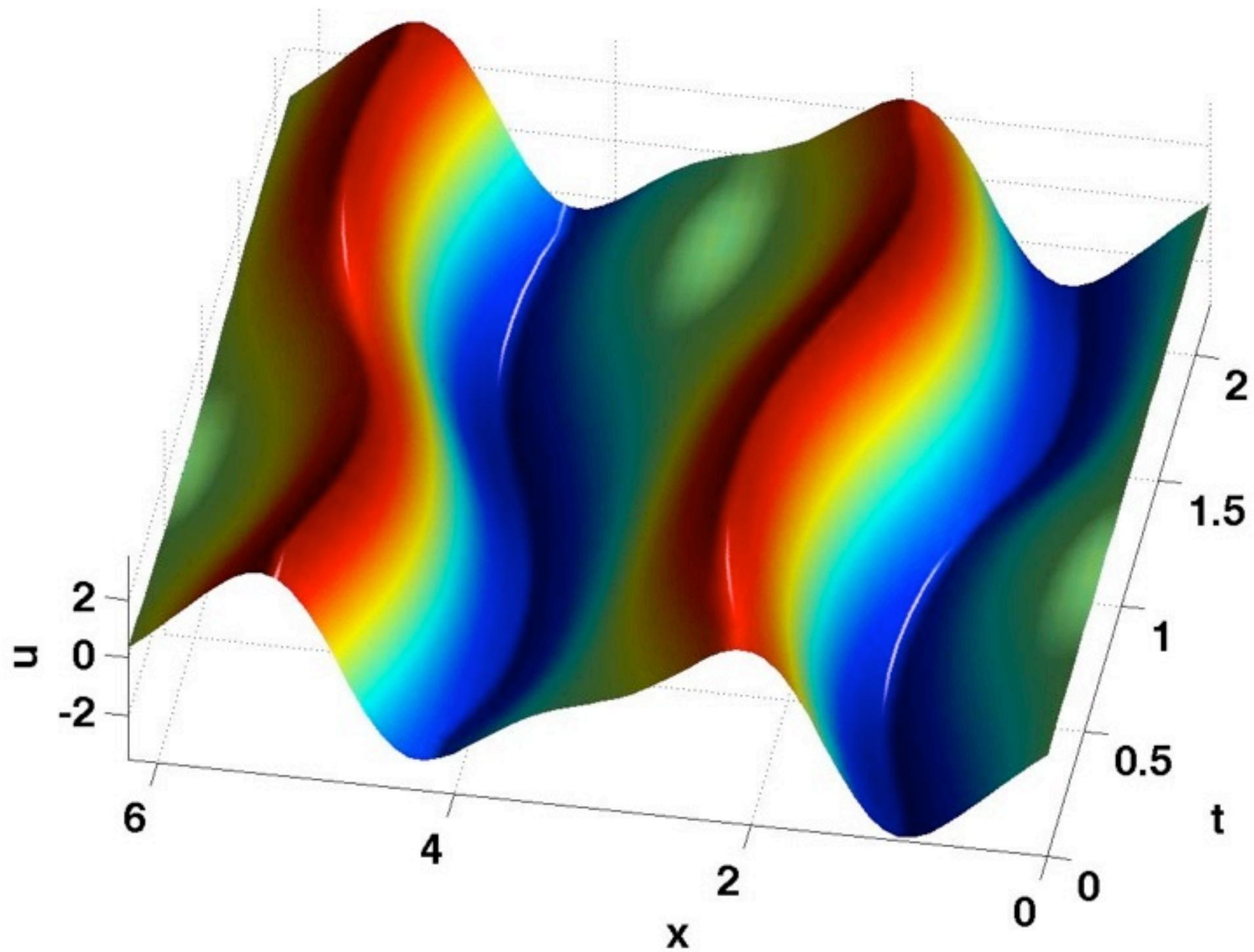
$$\nu = 0.12707$$



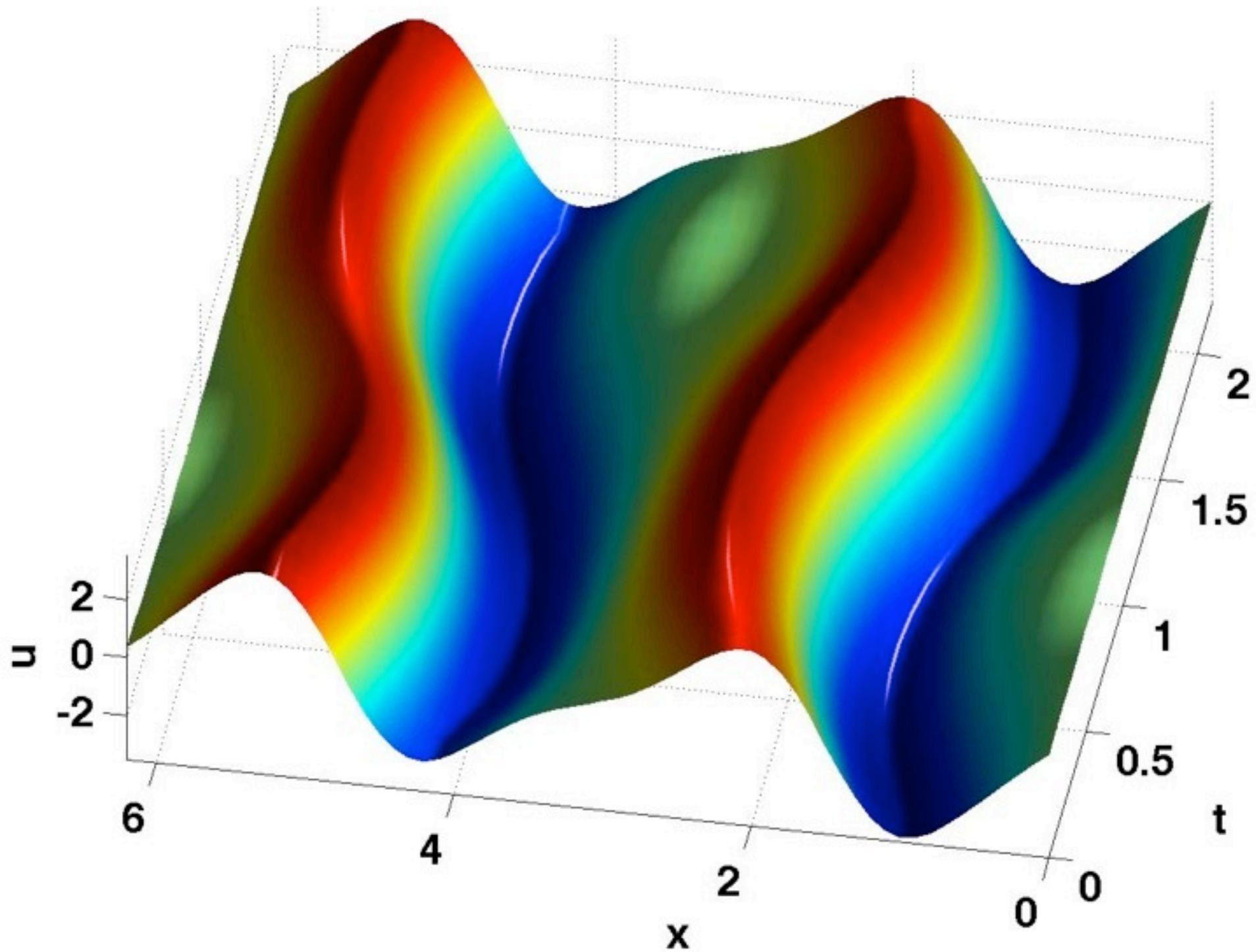
$$\nu = 0.12715$$



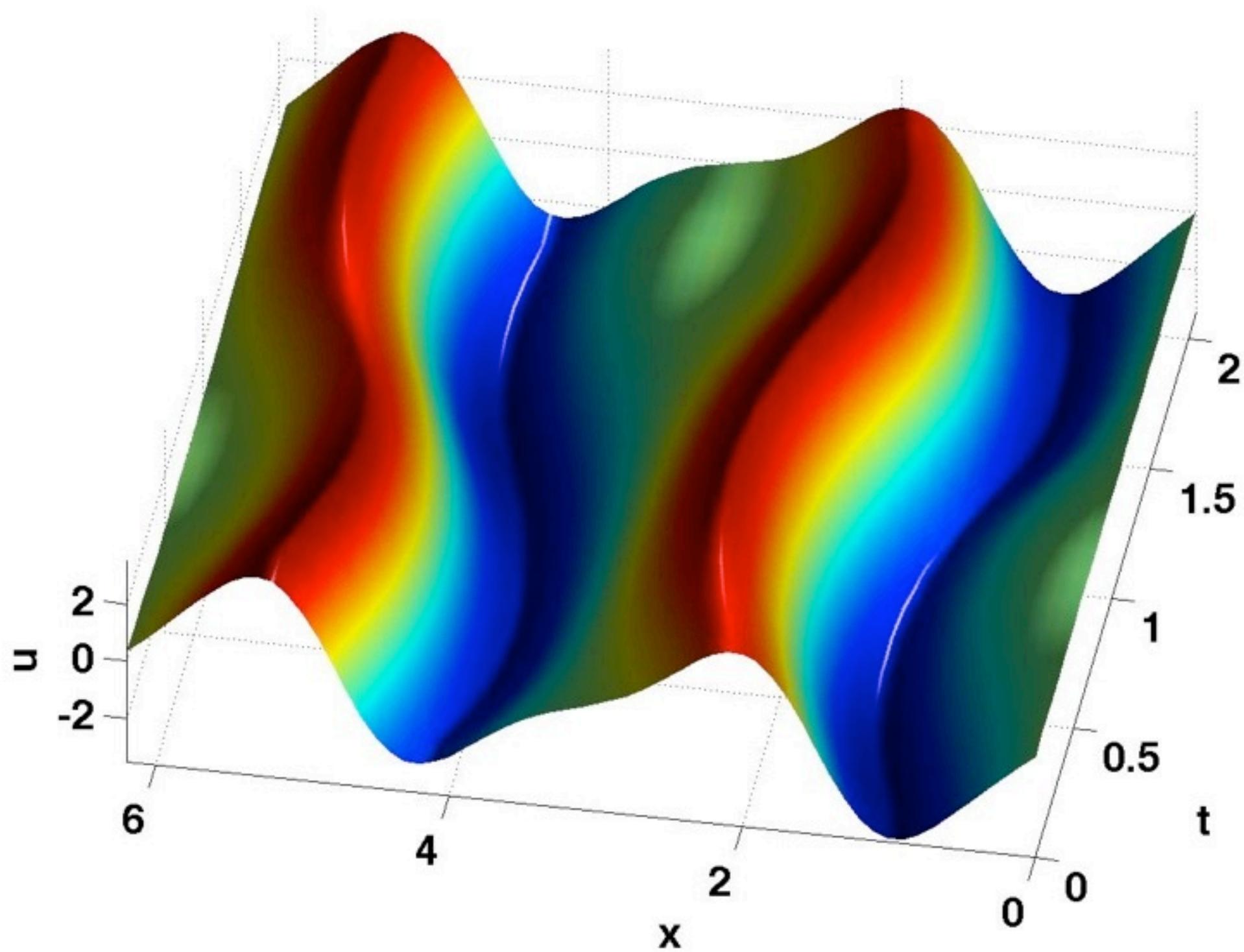
$$\nu = 0.12725$$



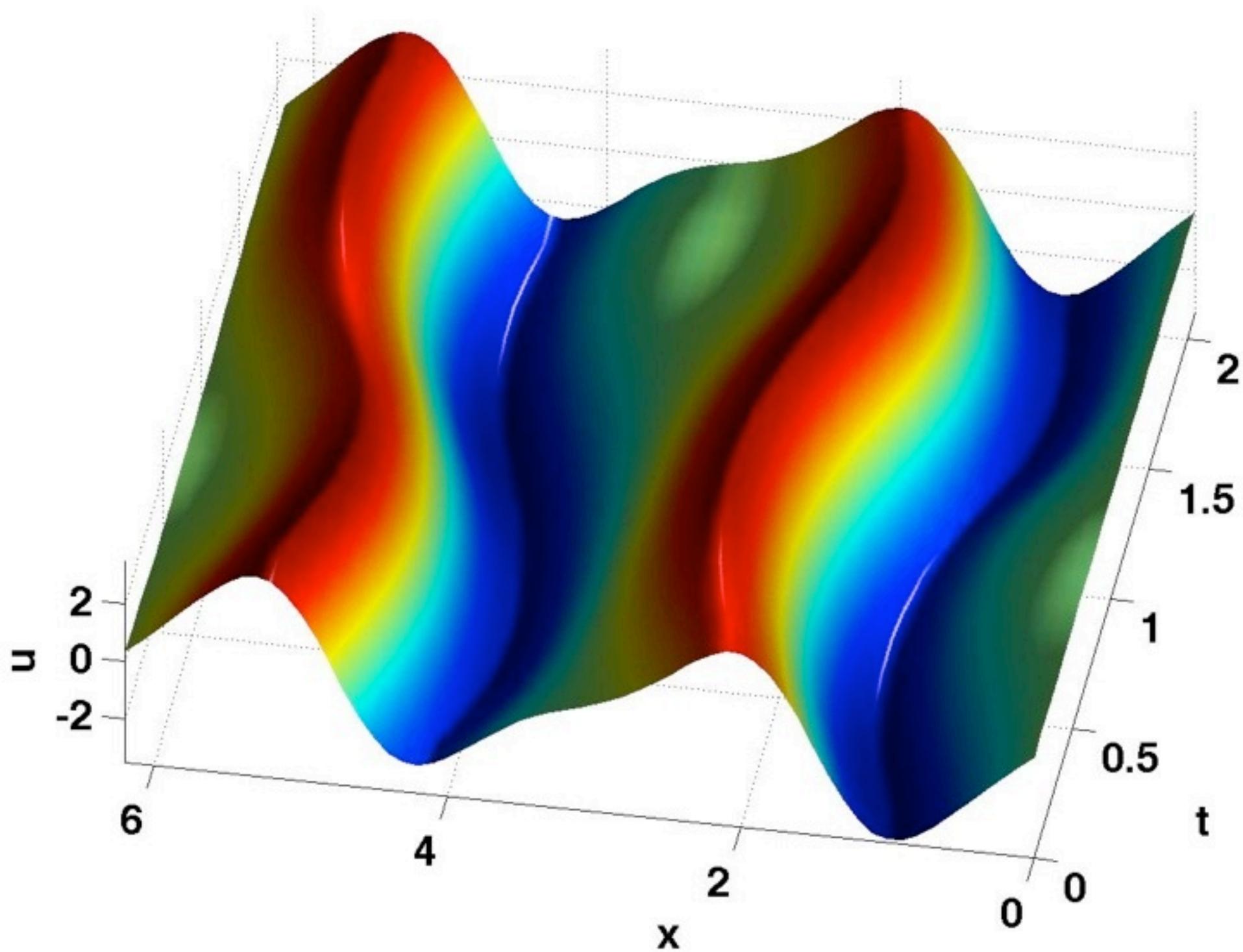
$$\nu = 0.12739$$



$$\nu = 0.12756$$



$$\nu = 0.12777$$



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